

**Imperial College
London**



Search for Hidden Particles

Ulrik Egede, on behalf of SHiP collaboration

UCL, 28 May 2014

A contradiction?

The Standard Model seems perfect and can exist without corrections to the Planck scale

But the same Standard Model can't explain dark matter, neutrino masses and baryogenesis.

I will propose a set of possible ways out of this

Outline

The physics landscape

Hidden sector theories

The neutrino minimal Standard Model

Design of a new beam-dump experiment

Sensitivity and future plans

Who are we?

www.cern.ch/ship

CERN-SPSC-2013-024 / SPSC-EOI-010

October 8, 2013

Proposal to Search for Heavy Neutral Leptons at the SPS

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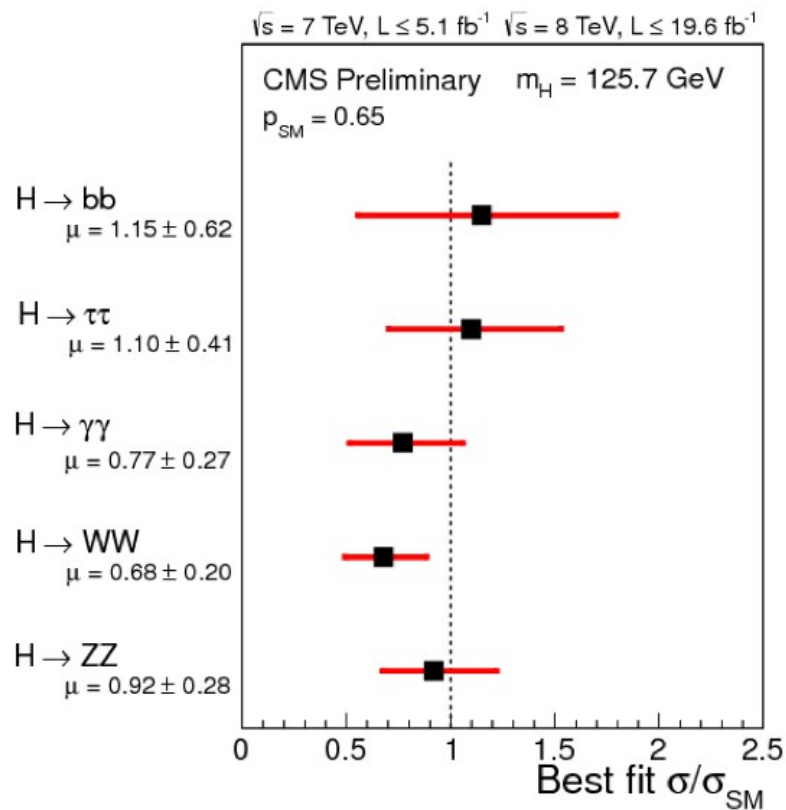
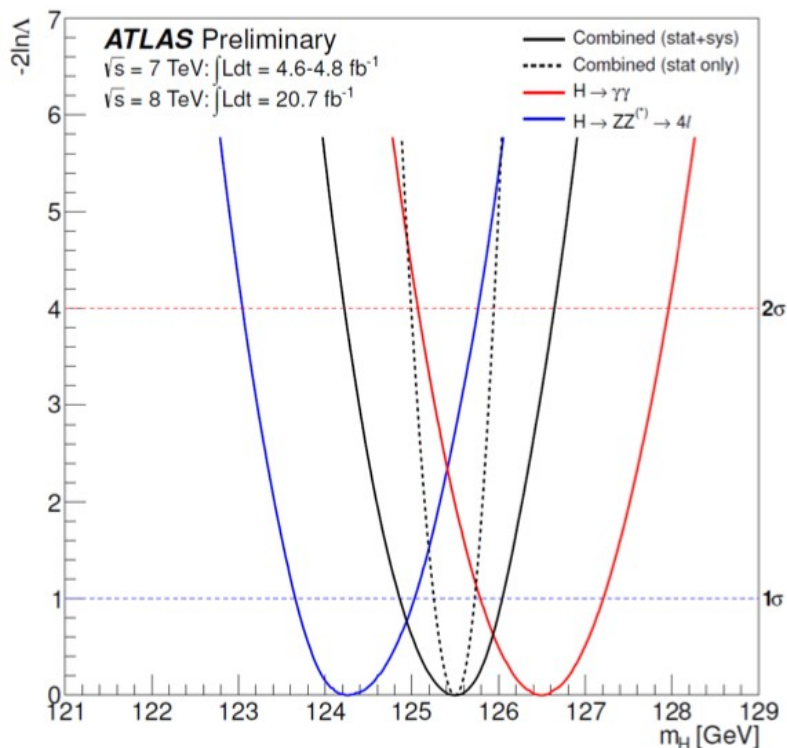
(†) *retired*

The triumph of the Standard Model

Boson consistent with the SM-Higgs has been found!

ATLAS : $M_H = 125.5 \pm 0.2$ (stat) $+0.5-0.6$ (syst) GeV/c^2

CMS : $M_H = 125.7 \pm 0.3$ (stat) ± 0.3 (syst) GeV/c^2



The triumph of the Standard Model

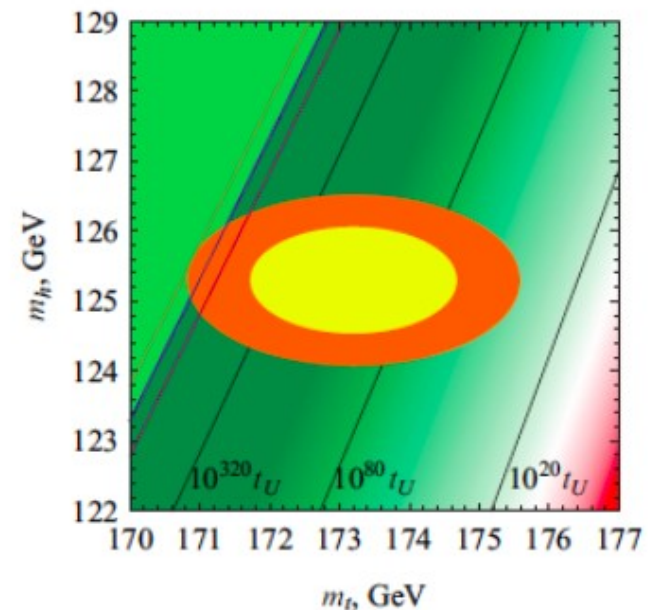
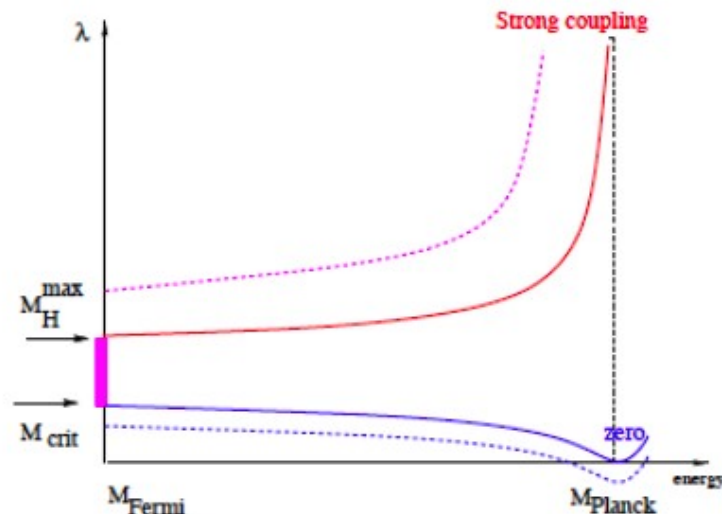
Mass value important for the stability of the vacuum:

$$M_H < 175 \text{ GeV}$$

SM weakly coupled up to the Plank energies !

$$M_H > 111 \text{ GeV}$$

EW vacuum is stable or metastable with a lifetime greatly exceeding the age of our Universe (JHEP 1208 (2012) 098)



The limitations of the Standard Model

But we still have a number of significant problems

Theory

- Radiative corrections to Higgs mass

 - fine-tuning

Experiment

- Matter anti-matter asymmetry in the Universe

- Neutrino masses and oscillations

- Non-baryonic dark matter

- Dark Energy

Direct searches ...

ATLAS SUSY Searches* - 95% CL Lower Limits

Status: Moriond 2014

ATLAS Preliminary

$$\int \mathcal{L} dt = (4.6 - 22.9) \text{ fb}^{-1} \quad \sqrt{s} = 7, 8 \text{ TeV}$$

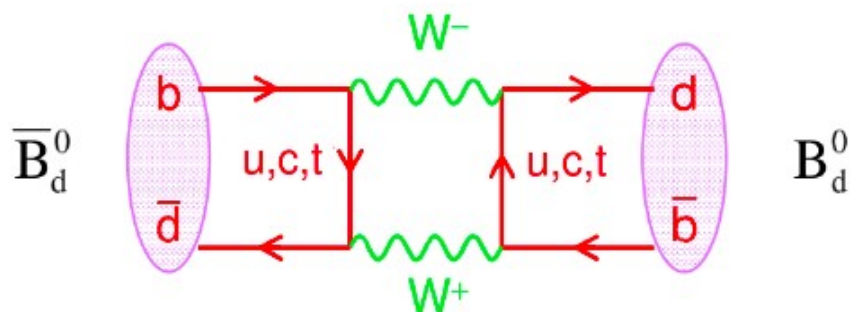
Model	e, μ, τ, γ	Jets	E_T^{miss}	$\int \mathcal{L} dt [\text{fb}^{-1}]$	Mass limit	Reference	
Inclusive Searches	MSUGRA/CMSSM	0	2-6 jets	Yes	20.3	\tilde{q}, \tilde{g} 1.7 TeV	$m(\tilde{q})=m(\tilde{g})$ ATLAS-CONF-2013-047
	MSUGRA/CMSSM	$1 e, \mu$	3-6 jets	Yes	20.3	\tilde{g} 1.2 TeV	any $m(\tilde{q})$ ATLAS-CONF-2013-062
	MSUGRA/CMSSM	0	7-10 jets	Yes	20.3	\tilde{g} 1.1 TeV	any $m(\tilde{q})$ 1308.1841
	$\tilde{q}\tilde{q}, \tilde{g} \rightarrow q\tilde{X}_1^0$	0	2-6 jets	Yes	20.3	\tilde{q} 740 GeV	$m(\tilde{X}_1^0)=0 \text{ GeV}$ ATLAS-CONF-2013-047
	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow q\tilde{X}_1^0$	0	2-6 jets	Yes	20.3	\tilde{g} 1.3 TeV	$m(\tilde{X}_1^0)=0 \text{ GeV}$ ATLAS-CONF-2013-047
	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow q\tilde{X}_1^0 \rightarrow qgW^{\pm}\tilde{X}_1^0$	$1 e, \mu$	3-6 jets	Yes	20.3	\tilde{g} 1.18 TeV	$m(\tilde{X}_1^0)<200 \text{ GeV}, m(\tilde{X}^{\pm})=0.5(m(\tilde{X}_1^0)+m(\tilde{g}))$ ATLAS-CONF-2013-062
	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow qg(\ell\ell)/(\nu\nu)\tilde{X}_1^0$	$2 e, \mu$	0-3 jets	-	20.3	\tilde{g} 1.12 TeV	$m(\tilde{X}_1^0)=0 \text{ GeV}$ ATLAS-CONF-2013-089
	GMSB ($\tilde{\ell}$ NLSP)	$2 e, \mu$	2-4 jets	Yes	4.7	\tilde{g} 1.24 TeV	$\tan\beta < 15$ 1208.4688
	GMSB ($\tilde{\ell}$ NLSP)	$1-2 \tau$	0-2 jets	Yes	20.7	\tilde{g} 1.4 TeV	$\tan\beta > 18$ ATLAS-CONF-2013-026
	GGM (bino NLSP)	2γ	-	Yes	20.3	\tilde{g} 1.28 TeV	$m(\tilde{X}_1^0)>50 \text{ GeV}$ ATLAS-CONF-2014-001
GGM (wino NLSP)	$1 e, \mu + \gamma$	-	Yes	4.8	\tilde{g} 619 GeV	$m(\tilde{X}_1^0)>50 \text{ GeV}$ ATLAS-CONF-2012-144	
GGM (higgsino-bino NLSP)	γ	1 b	Yes	4.8	\tilde{g} 900 GeV	$m(\tilde{X}_1^0)>220 \text{ GeV}$ 1211.1167	
GGM (higgsino NLSP)	$2 e, \mu (Z)$	0-3 jets	Yes	5.8	\tilde{g} 690 GeV	$m(H)>200 \text{ GeV}$ ATLAS-CONF-2012-152	
Gravitino LSP	0	mono-jet	Yes	10.5	\tilde{g}/\tilde{X}_1^0 scale 645 GeV	$m(\tilde{g})>10^{-1} \text{ eV}$ ATLAS-CONF-2012-147	
3 rd gen. \tilde{g} med.	$\tilde{g} \rightarrow b\tilde{X}_1^0$	0	3 b	Yes	20.1	\tilde{g} 1.2 TeV	$m(\tilde{X}_1^0)<600 \text{ GeV}$ ATLAS-CONF-2013-061
	$\tilde{g} \rightarrow t\tilde{X}_1^0$	0	7-10 jets	Yes	20.3	\tilde{g} 1.1 TeV	$m(\tilde{X}_1^0) < 350 \text{ GeV}$ 1308.1841
	$\tilde{g} \rightarrow t\tilde{X}_1^0$	$0-1 e, \mu$	3 b	Yes	20.1	\tilde{g} 1.34 TeV	$m(\tilde{X}_1^0)<400 \text{ GeV}$ ATLAS-CONF-2013-061
	$\tilde{g} \rightarrow b\tilde{X}_1^0$	$0-1 e, \mu$	3 b	Yes	20.1	\tilde{g} 1.3 TeV	$m(\tilde{X}_1^0)<300 \text{ GeV}$ ATLAS-CONF-2013-061
3 rd gen. squarks direct production	$\tilde{b}_1\tilde{b}_1, \tilde{b}_1 \rightarrow b\tilde{X}_1^0$	0	2 b	Yes	20.1	\tilde{b}_1 100-620 GeV	$m(\tilde{X}_1^0)<90 \text{ GeV}$ 1308.2631
	$\tilde{b}_1\tilde{b}_1, \tilde{b}_1 \rightarrow t\tilde{X}_1^0$	$2 e, \mu (SS)$	0-3 b	Yes	20.7	\tilde{b}_1 275-430 GeV	$m(\tilde{X}_1^0)=2 m(\tilde{X}_1^0)$ ATLAS-CONF-2013-007
	$\tilde{t}_1\tilde{t}_1$ (light), $\tilde{t}_1 \rightarrow b\tilde{X}_1^0$	$1-2 e, \mu$	1-2 b	Yes	4.7	\tilde{t}_1 110-167 GeV	$m(\tilde{X}_1^0)=55 \text{ GeV}$ 1208.4305, 1209.2102
	$\tilde{t}_1\tilde{t}_1$ (light), $\tilde{t}_1 \rightarrow Wb\tilde{X}_1^0$	$2 e, \mu$	0-2 jets	Yes	20.3	\tilde{t}_1 130-210 GeV	$m(\tilde{X}_1^0)=m(\tilde{t}_1)-m(W)-50 \text{ GeV}, m(\tilde{t}_1)<m(\tilde{X}_1^0)$ 1403.4853
	$\tilde{t}_1\tilde{t}_1$ (medium), $\tilde{t}_1 \rightarrow t\tilde{X}_1^0$	$2 e, \mu$	2 jets	Yes	20.3	\tilde{t}_1 215-530 GeV	$m(\tilde{X}_1^0)=1 \text{ GeV}$ 1403.4853
	$\tilde{t}_1\tilde{t}_1$ (medium), $\tilde{t}_1 \rightarrow b\tilde{X}_1^0$	0	2 b	Yes	20.1	\tilde{t}_1 150-580 GeV	$m(\tilde{X}_1^0)<200 \text{ GeV}, m(\tilde{X}_1^0)-m(\tilde{X}_1^0)=5 \text{ GeV}$ 1308.2631
	$\tilde{t}_1\tilde{t}_1$ (heavy), $\tilde{t}_1 \rightarrow t\tilde{X}_1^0$	$1 e, \mu$	1 b	Yes	20.7	\tilde{t}_1 200-610 GeV	$m(\tilde{X}_1^0)=0 \text{ GeV}$ ATLAS-CONF-2013-037
	$\tilde{t}_1\tilde{t}_1$ (heavy), $\tilde{t}_1 \rightarrow b\tilde{X}_1^0$	0	2 b	Yes	20.5	\tilde{t}_1 320-660 GeV	$m(\tilde{X}_1^0)=0 \text{ GeV}$ ATLAS-CONF-2013-024
	$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow c\tilde{X}_1^0$	0	mono-jet/c-tag	Yes	20.3	\tilde{t}_1 90-200 GeV	$m(\tilde{t}_1)-m(\tilde{X}_1^0)<85 \text{ GeV}$ ATLAS-CONF-2013-068
	$\tilde{t}_1\tilde{t}_1$ (natural GMSB)	$2 e, \mu (Z)$	1 b	Yes	20.3	\tilde{t}_1 150-580 GeV	$m(\tilde{X}_1^0)>150 \text{ GeV}$ 1403.5222
$\tilde{t}_2\tilde{t}_2, \tilde{t}_2 \rightarrow \tilde{t}_1 + Z$	$3 e, \mu (Z)$	1 b	Yes	20.3	\tilde{t}_2 290-600 GeV	$m(\tilde{X}_1^0)<200 \text{ GeV}$ 1403.5222	
EW direct	$\tilde{\ell}_L\tilde{\ell}_L, \tilde{\ell} \rightarrow \ell\tilde{X}_1^0$	$2 e, \mu$	0	Yes	20.3	$\tilde{\ell}$ 90-325 GeV	$m(\tilde{X}_1^0)=0 \text{ GeV}$ 1403.5294
	$\tilde{\ell}_L\tilde{\ell}_L, \tilde{\ell} \rightarrow \ell(\tilde{\nu})$	$2 e, \mu$	0	Yes	20.3	$\tilde{\ell}_L^{\pm}$ 140-465 GeV	$m(\tilde{X}_1^0)=0 \text{ GeV}, m(\tilde{\ell}, \tilde{\nu})=0.5(m(\tilde{X}_1^0)+m(\tilde{\ell}_L^{\pm}))$ 1403.5294
	$\tilde{\ell}_L\tilde{\ell}_L, \tilde{\ell} \rightarrow \tilde{\nu}(\tau\tilde{\nu})$	2τ	-	Yes	20.7	$\tilde{\ell}_L^{\pm}$ 180-330 GeV	$m(\tilde{X}_1^0)=0 \text{ GeV}, m(\tilde{\ell}, \tilde{\nu})=0.5(m(\tilde{X}_1^0)+m(\tilde{\ell}_L^{\pm}))$ ATLAS-CONF-2013-028
	$\tilde{\ell}_L\tilde{\ell}_L \rightarrow \tilde{\ell}_L \nu(\tilde{\ell}(\tilde{\nu}\nu), \tilde{\ell}\tilde{\nu}(\tilde{\ell}(\tilde{\nu}\nu)))$	$3 e, \mu$	0	Yes	20.3	$\tilde{\ell}_L^{\pm}, \tilde{\ell}_L^0$ 700 GeV	$m(\tilde{X}_1^0)=m(\tilde{X}_2^0), m(\tilde{X}_1^0)=0, m(\tilde{\ell}, \tilde{\nu})=0.5(m(\tilde{X}_1^0)+m(\tilde{X}_2^0))$ 1402.7029
	$\tilde{\ell}_L\tilde{\ell}_L \rightarrow W\tilde{\ell}_L^0 Z\tilde{X}_1^0$	$2-3 e, \mu$	0	Yes	20.3	$\tilde{\ell}_L^{\pm}, \tilde{\ell}_L^0$ 420 GeV	$m(\tilde{X}_1^0)=m(\tilde{X}_2^0), m(\tilde{X}_1^0)=0, \text{ sleptons decoupled}$ 1403.5294, 1402.7029
	$\tilde{\ell}_L\tilde{\ell}_L \rightarrow W\tilde{\ell}_L^0 h\tilde{X}_1^0$	$1 e, \mu$	2 b	Yes	20.3	$\tilde{\ell}_L^{\pm}, \tilde{\ell}_L^0$ 285 GeV	$m(\tilde{X}_1^0)=m(\tilde{X}_2^0), m(\tilde{X}_1^0)=0, \text{ sleptons decoupled}$ ATLAS-CONF-2013-093
	$\tilde{\ell}_L\tilde{\ell}_L, \tilde{\ell} \rightarrow \tilde{\ell}_L \tilde{X}_1^0$	-	-	-	-	$\tilde{\ell}_L^{\pm}, \tilde{\ell}_L^0$ 270 GeV	$m(\tilde{X}_1^0)-m(\tilde{X}_2^0)=160 \text{ MeV}, \tau(\tilde{X}_1^0)=0.2 \text{ ns}$ ATLAS-CONF-2013-069
Long-lived particles	Stable, stopped \tilde{g} R-hadron	0	1-5 jets	Yes	22.9	\tilde{g} 832 GeV	$m(\tilde{X}_1^0)=100 \text{ GeV}, 10 \mu\text{s} < \tau(\tilde{g}) < 1000 \text{ s}$ ATLAS-CONF-2013-057
	GMSB, stable $\tilde{\tau}, \tilde{\tau} \rightarrow \tilde{\tau}(\tilde{\nu}, \tilde{\mu}) + \tau(e, \mu)$	$1-2 \mu$	-	-	15.9	$\tilde{\tau}$ 475 GeV	$10 < \tan\beta < 50$ ATLAS-CONF-2013-058
	GMSB, $\tilde{\chi}_1^0 \rightarrow \gamma\tilde{Z}, \text{ long-lived } \tilde{\chi}_1^0$	2γ	-	Yes	4.7	$\tilde{\chi}_1^0$ 230 GeV	$0.4 < \tau(\tilde{\chi}_1^0) < 2 \text{ ns}$ 1304.6310
	$\tilde{q}\tilde{q}, \tilde{\chi}_1^0 \rightarrow q\tilde{q}\mu$ (RPV)	$1 \mu, \text{ displ. vtx}$	-	-	20.3	\tilde{q} 1.0 TeV	$1.5 < c\tau < 156 \text{ mm}, \text{BR}(\mu)=1, m(\tilde{X}_1^0)=108 \text{ GeV}$ ATLAS-CONF-2013-092
	LFV $pp \rightarrow \tilde{\nu}_e + X, \tilde{\nu}_e \rightarrow e + \mu$	$2 e, \mu$	-	-	4.6	$\tilde{\nu}_e$ 1.61 TeV	$\lambda'_{111}=0.10, \lambda'_{132}=0.05$ 1212.1272
	LFV $pp \rightarrow \tilde{\nu}_e + X, \tilde{\nu}_e \rightarrow e(\mu) + \tau$	$1 e, \mu + \tau$	-	-	4.6	$\tilde{\nu}_e$ 1.1 TeV	$\lambda'_{111}=0.10, \lambda'_{12133}=0.05$ 1212.1272
	Bi-linear RPV CMSSM	$1 e, \mu$	7 jets	Yes	4.7	\tilde{q}, \tilde{g} 1.2 TeV	$m(\tilde{q})=m(\tilde{g}), c\tau_{LSP} < 1 \text{ mm}$ ATLAS-CONF-2012-140
$\tilde{\chi}_1^0\tilde{\chi}_1^0, \tilde{\chi}_1^0 \rightarrow W\tilde{X}_1^0, \tilde{\chi}_1^0 \rightarrow e\tilde{\nu}_\mu, e\mu\tilde{\nu}_e$	$4 e, \mu$	-	Yes	20.7	$\tilde{\chi}_1^0$ 760 GeV	$m(\tilde{X}_1^0)>300 \text{ GeV}, \lambda'_{121}>0$ ATLAS-CONF-2013-036	
$\tilde{\chi}_1^0\tilde{\chi}_1^0, \tilde{\chi}_1^0 \rightarrow W\tilde{X}_1^0, \tilde{\chi}_1^0 \rightarrow \tau\tilde{\nu}_e, e\tau\tilde{\nu}_e$	$3 e, \mu + \tau$	-	Yes	20.7	$\tilde{\chi}_1^0$ 350 GeV	$m(\tilde{X}_1^0)>80 \text{ GeV}, \lambda'_{133}>0$ ATLAS-CONF-2013-036	
$\tilde{g} \rightarrow q\tilde{q}q$	0	6-7 jets	-	20.3	\tilde{g} 916 GeV	$\text{BR}(\tau)=\text{BR}(b)=\text{BR}(c)=0\%$ ATLAS-CONF-2013-091	
$\tilde{g} \rightarrow \tilde{t}_1 t, \tilde{t}_1 \rightarrow b\tilde{s}$	$2 e, \mu (SS)$	0-3 b	Yes	20.7	\tilde{g} 880 GeV	ATLAS-CONF-2013-007	
Other	Scalar gluon pair, $sgluon \rightarrow q\tilde{q}$	0	4 jets	-	4.6	$sgluon$ 100-287 GeV	incl. limit from 1110.2693 1210.4826
	Scalar gluon pair, $sgluon \rightarrow t\tilde{t}$	$2 e, \mu (SS)$	2 b	Yes	14.3	$sgluon$ 350-800 GeV	ATLAS-CONF-2013-051
	WIMP interaction (D5, Dirac χ)	0	mono-jet	Yes	10.5	\tilde{M}^* scale 704 GeV	$m(\chi)<80 \text{ GeV}, \text{ limit of } < 687 \text{ GeV for D8}$ ATLAS-CONF-2012-147

$\sqrt{s} = 7 \text{ TeV}$ full data
 $\sqrt{s} = 8 \text{ TeV}$ partial data
 $\sqrt{s} = 8 \text{ TeV}$ full data

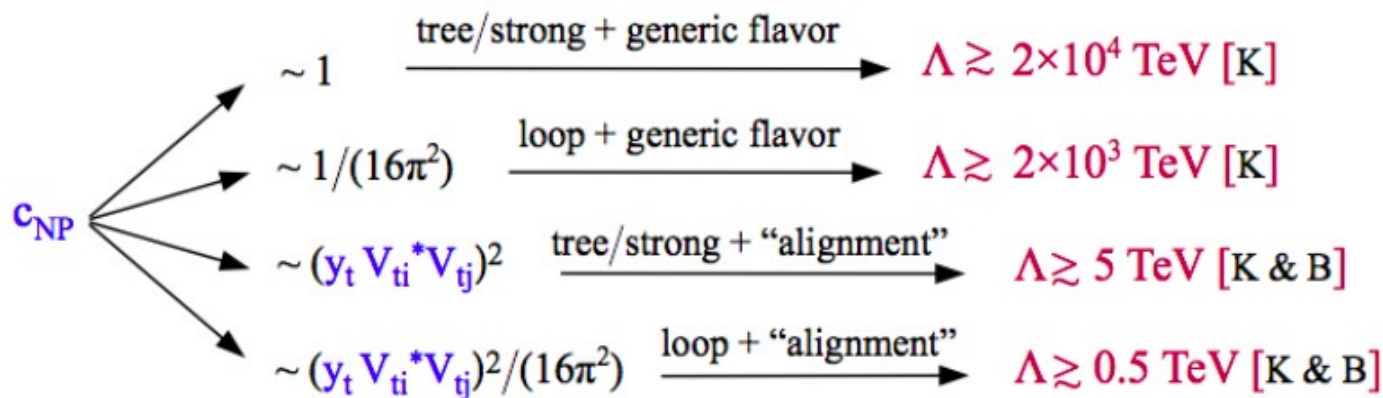
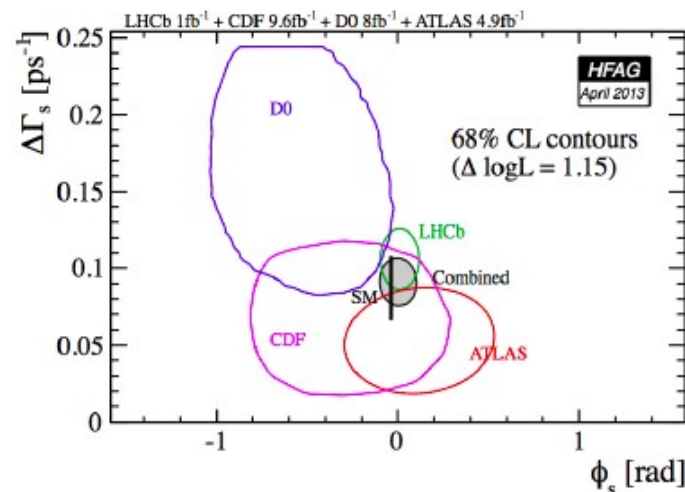
Mass scale [TeV]

*Only a selection of the available mass limits on new states or phenomena is shown. All limits quoted are observed minus 1σ theoretical signal cross section uncertainty.

... and indirect searches



$$M(B_d^0 - \bar{B}_d^0) \sim \frac{(y_t^2 V_{tb}^* V_{td})^2}{16\pi^2 m_t^2} + c_{NP} \frac{1}{\Lambda^2}$$



A hidden sector?

Rather than being heavy, could new particles be light but very weakly interacting?

A new, light “hidden sector” of particles

Singlets with respect to gauge group of the SM

How could we have missed it?

Key is that it only interacts with SM particles through some kind of mixing through a “portal” particle

A hidden sector?

Several possibilities for renormalisable singlet operators

Vector portal, U(1) $B_{\mu\nu}$

Massive vector photon (paraphoton, secluded photon...)
mixing with regular photon $\rightarrow \epsilon B_{\mu\nu} F^{\mu\nu}$

Higgs portal

Scalar field χ , $(\mu\chi + \lambda\chi^2)H^\dagger H$

Axial portal

Pseudo Nambu-Goldstone bosons

Axion like vector field , $(a/F)G_{\mu\nu} G^{\mu\nu}$, $(\partial_\mu a/F)\psi^\dagger \gamma_\mu \gamma_5 \psi$

Neutrino portal

Heavy neutral leptons (HNL), $YH^T N' L$

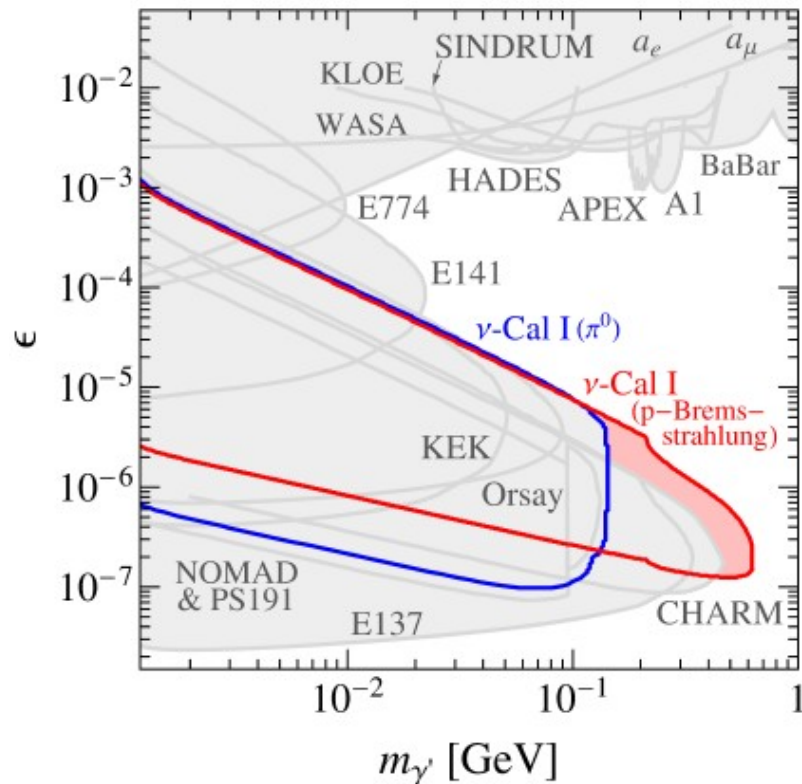
Low energy SYSY

Vector Portal

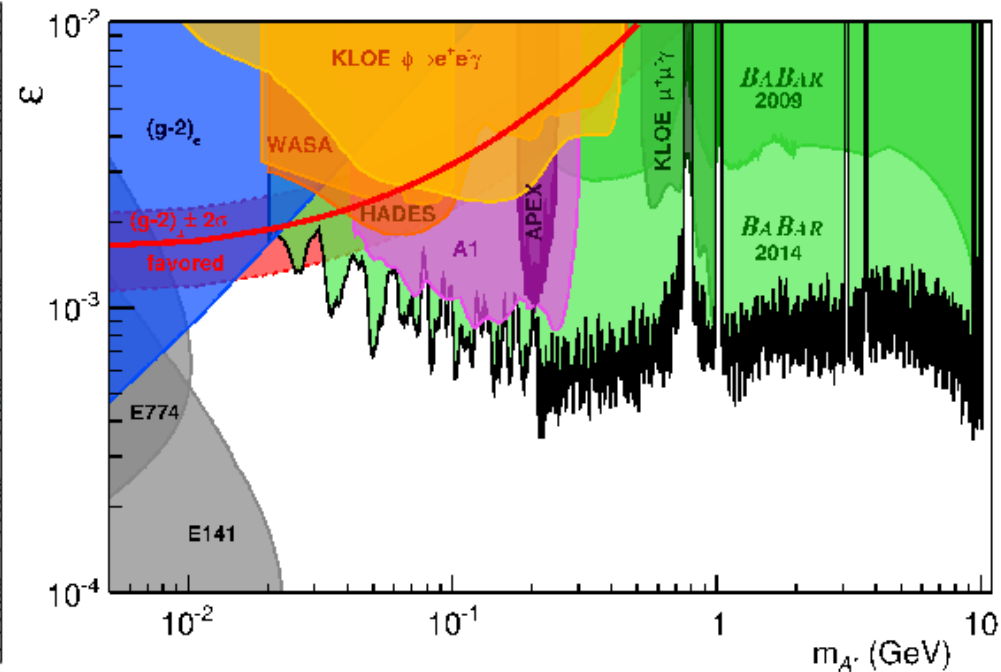
Exploit mixing between a virtual photon and the dark photon

No other interactions with SM particles - “light-shining-through-a-wall” experiments

[arXiv:1311.3870]



[arXiv:1406.2980]



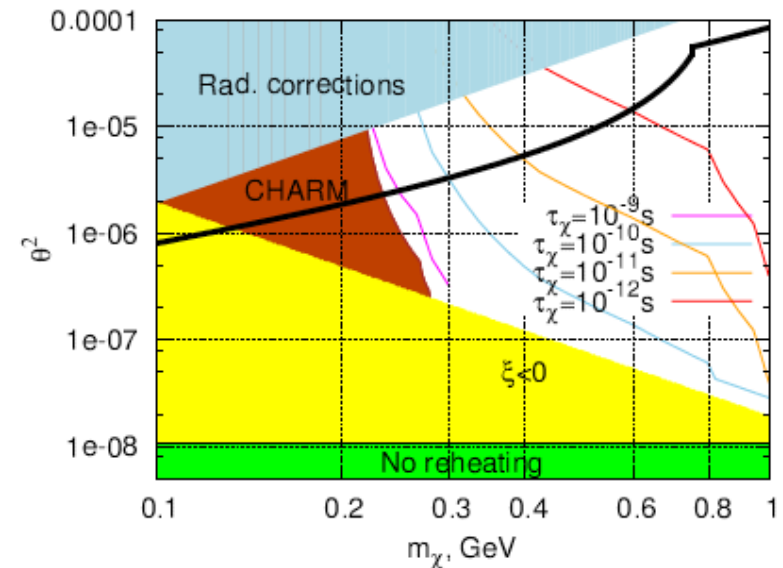
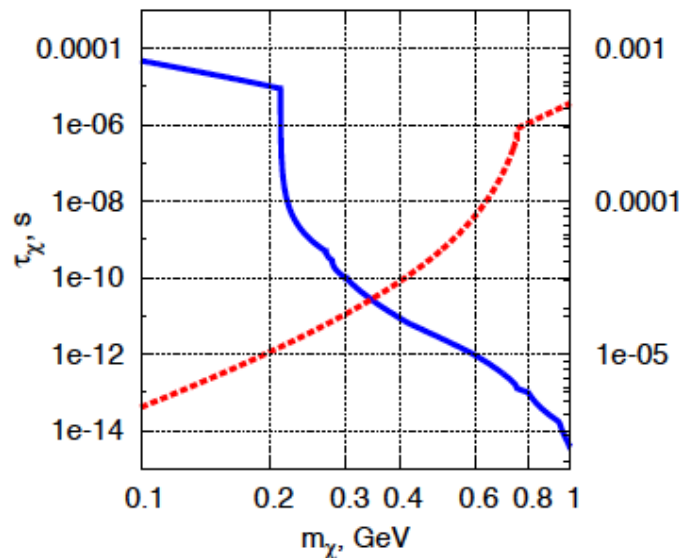
Higgs Portal

Example of inflaton

Together with Higgs, generates inflation of the early Universe
 Model has a 7 keV (warm) DM candidate and respects constraints from BICEP2 and Planck

Interesting mass region $0.3 \text{ GeV} < m_\chi < 1 \text{ GeV}$

Little experimental exploration of interesting region...



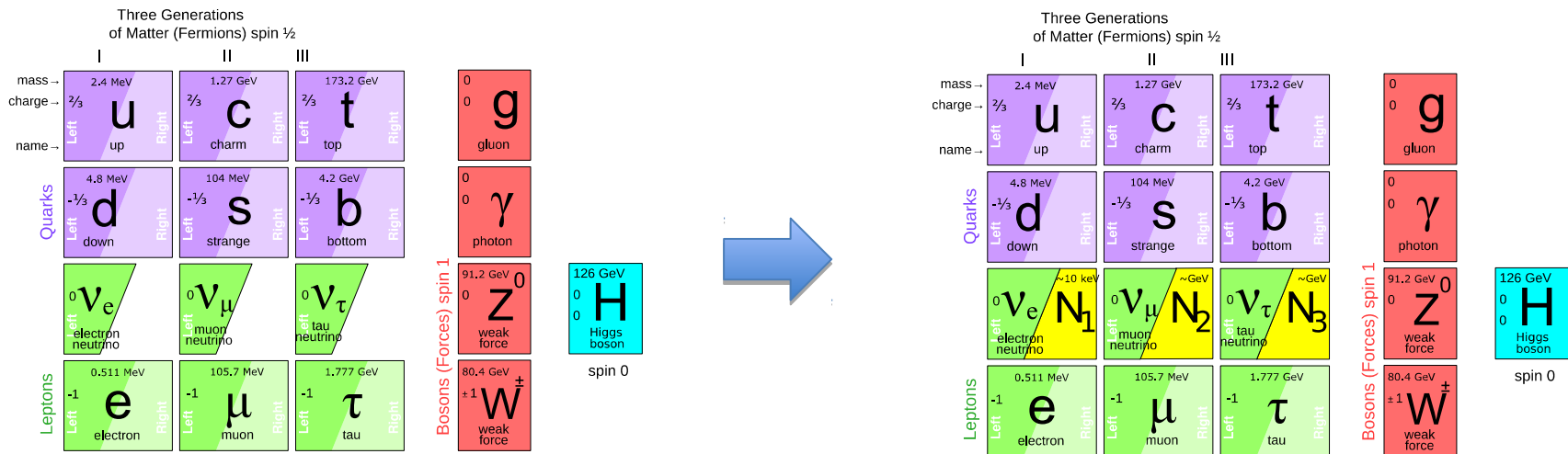
Neutrino Portal

[T.Asaka, M.Shaposhnikov, Phys. Lett B620 (2005) 17

The neutrino Minimal Standard Model (ν MSSM) aims to explain

Matter anti-matter asymmetry in the Universe, neutrino masses and oscillations, non-baryonic dark matter

Adds three right-handed, Majorana, Heavy Neutral Leptons (HNL), N_1 , N_2 and N_3



Neutrino Portal

[T.Asaka, M.Shaposhnikov, Phys. Lett B620 (2005) 17

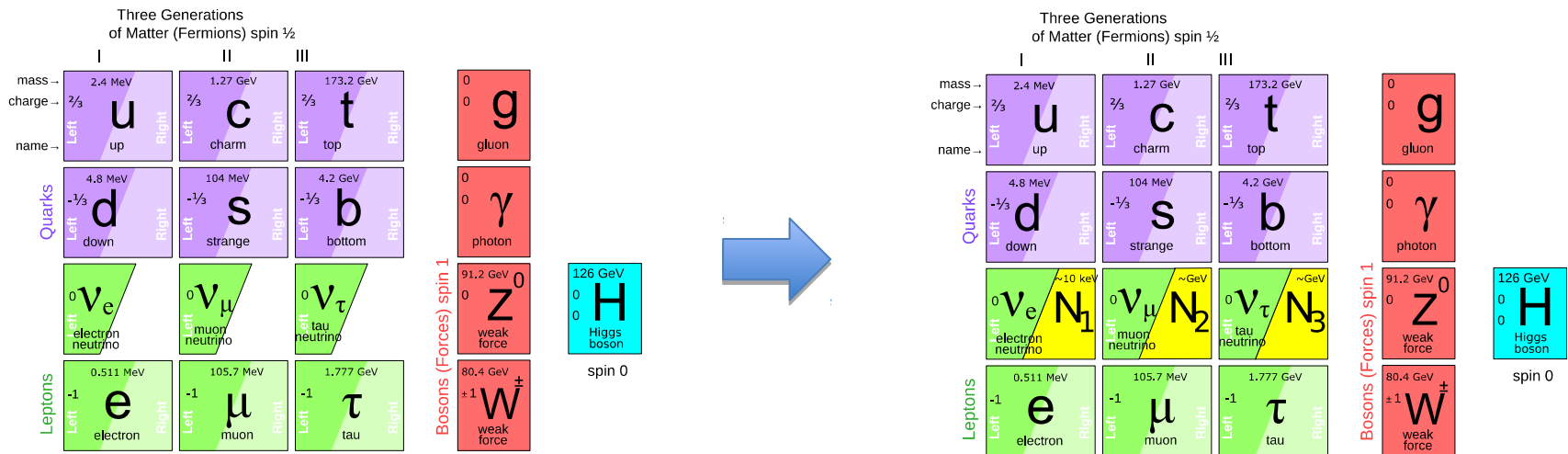
N_1

Mass in keV region, (warm) dark matter candidate

$N_{2,3}$

Mass in 100 MeV – GeV region

Generate neutrino masses via see-saw and produce baryon asymmetry of the Universe



See-saw for ν mass

Most general renormalisable Lagrangian of all SM particles (+3 singlets wrt the SM gauge group):

$$L_{\text{singlet}} = i\bar{N}_I \partial_\mu \gamma^\mu N_I - Y_{I\alpha} \bar{N}_I^c \tilde{H} L_\alpha^c - M_I \bar{N}_I^c N_I + \text{h.c.},$$

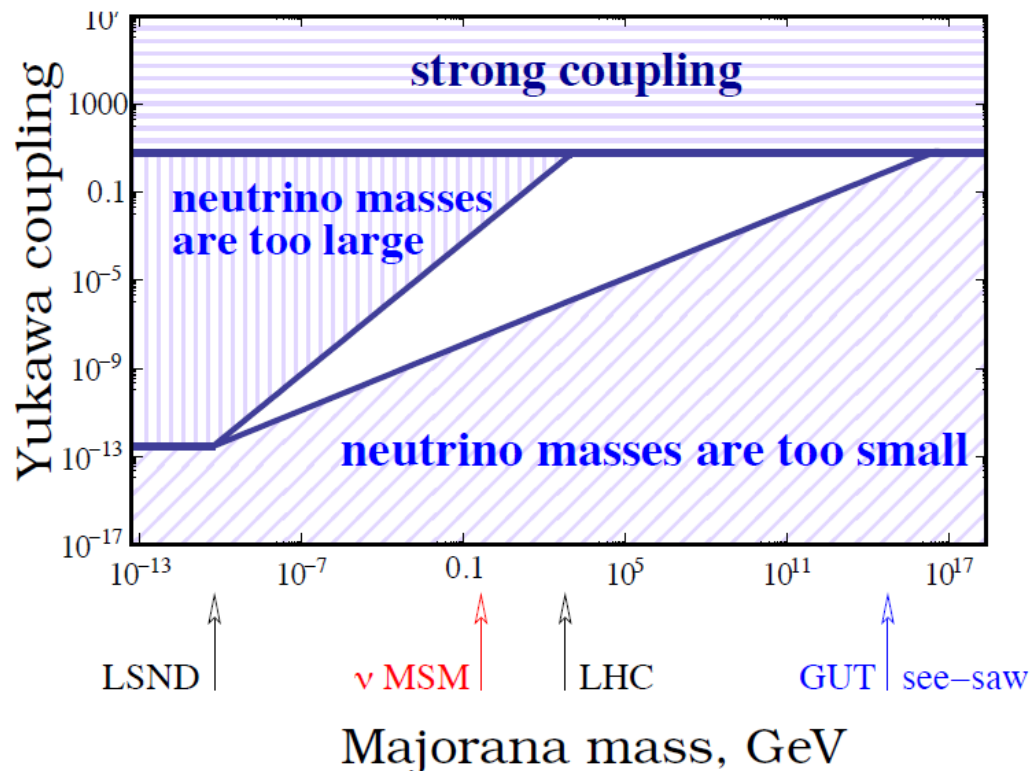
Yukawa term: mixing of N_I with active neutrinos to explain oscillations

Majorana term which carries no gauge charge

See-saw for ν mass

The scale of the active neutrino mass is given by the see-saw formula, $m_\nu = m_D^2/M$

Typical value of the Dirac mass term is linked to the Yukawa coupling of the l -th neutrino by $m_D \sim Y_{la} v$



Constraints on N_1 as DM

Stability

Must have a lifetime larger than that of the Universe

Production

Created in the early Universe in reactions $l^+l^- \rightarrow \nu N_1$,
 $qq \rightarrow \nu N_1$ etc.

Need to provide correct DM abundance

Structure formation

Should be heavy enough to not erase non-uniformities at small scales

Decay

Should not produce decays we have already excluded!

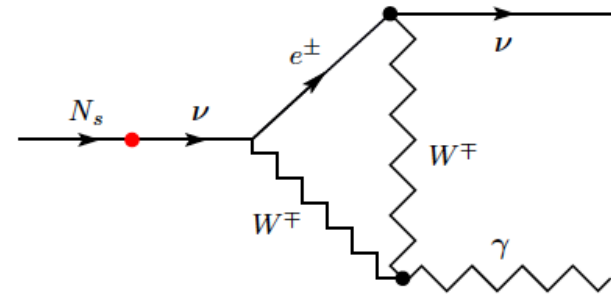
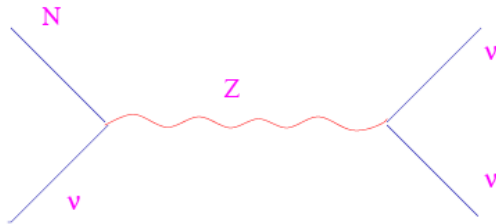
N_1 – dark matter candidate

Small Yukawa couplings mean that N_1 can be very stable

$$\tau_{N_1} = 10^{14} \text{ years} \left(\frac{10 \text{ keV}}{M_N} \right)^5 \left(\frac{10^{-8}}{\theta_1^2} \right) \quad \theta_1 = \frac{m_D}{M_N}$$

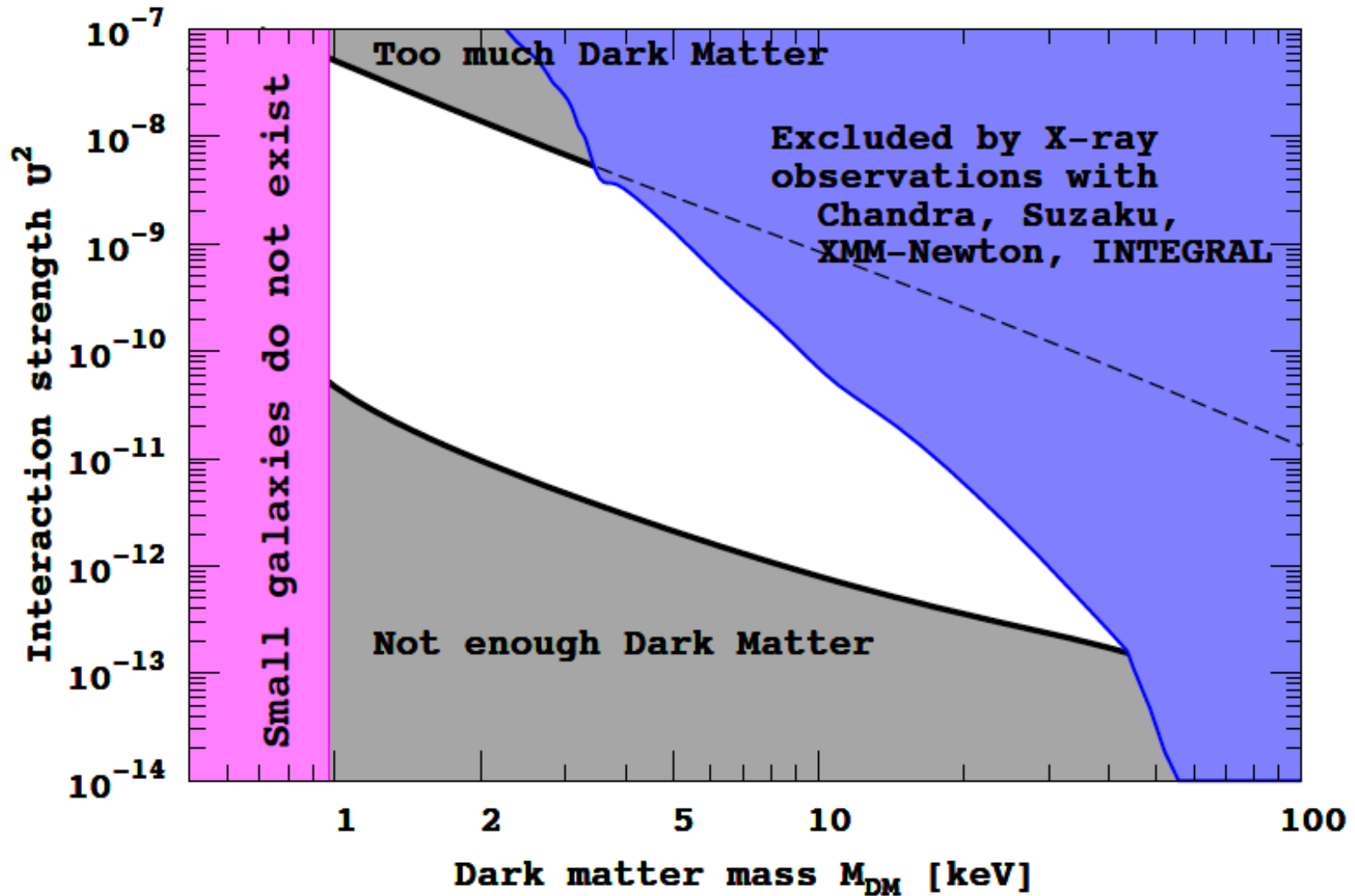
Main decay mode $N \rightarrow 3\nu$, clearly unobservable

Subdominant radiative decay $N \rightarrow \nu\gamma$ would give a monoenergetic photon with $E_\gamma = M_N/2$



N_1 allowed parameter space

20



New line in galaxy spectrum?

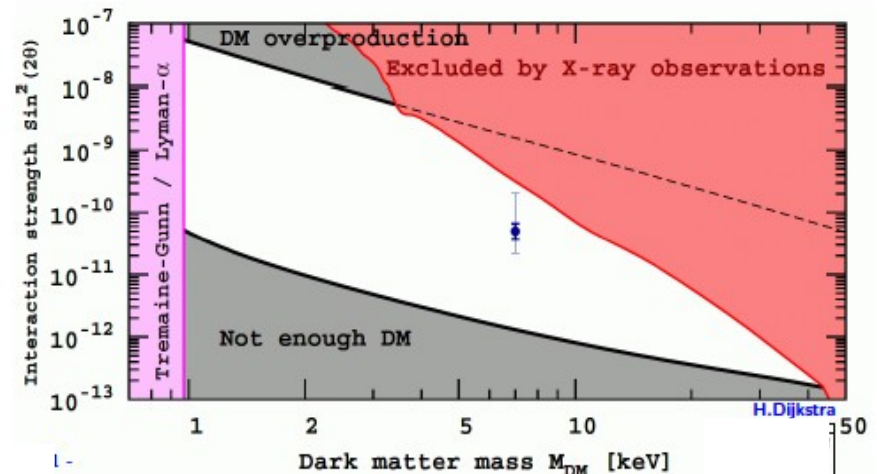
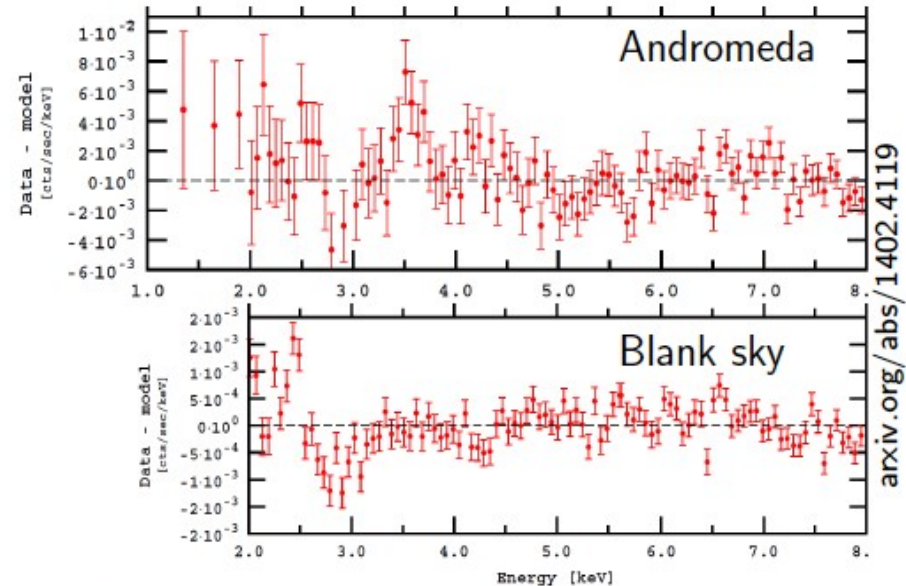
ArXiv:1402.4119

An unidentified line in the x-ray spectrum of the Andromeda galaxy and Perseus galaxy cluster $E_{\gamma} \sim 3.5$ keV

ArXiv:1402.2301

Detection of an unidentified emission line in the stacked x-ray spectrum of galaxy clusters $E_{\gamma} \sim 3.56$ keV

Astro-H will be able to check these claims with better energy resolution



Generating the baryon asymmetry

CP is not conserved in the ν MSM :

6 new CP-violating phases in lepton sector

Process for Baryon asymmetry

HNL are created in the early Universe

CPV in the interference of HNL production and decay

Lepton number asymmetry goes from HNL to active neutrinos

Asymmetry transferred to baryons via “sphaleron processes”

$N_{2,3}$ production and decay

$M(N_2) \approx M(N_3) \sim$ a few GeV \rightarrow can dramatically increase amount of CPV to explain Baryon Asymmetry of the Universe (BAU)

Explanation of DM with N_1 reduces number of free parameters, need degeneracy to ensure sufficient CPV

Very weak $N_{2,3}$ to ν mixing ($\sim U^2$) \rightarrow $N_{2,3}$ are much longer-lived than the SM particles

Typical lifetimes >10 ms for $(N_{2,3}) \sim 1$ GeV \rightarrow decay distance O(km)

Too large U erases any BAU

$$\tau = \frac{U^2 G_F^2 M_N^5}{86 \pi^3}$$

$N_{2,3}$ production and decay

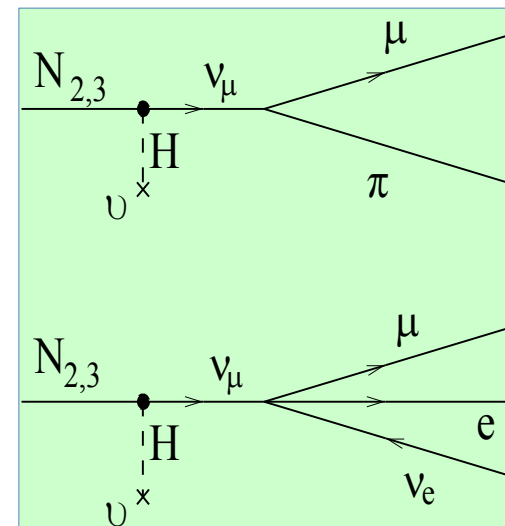
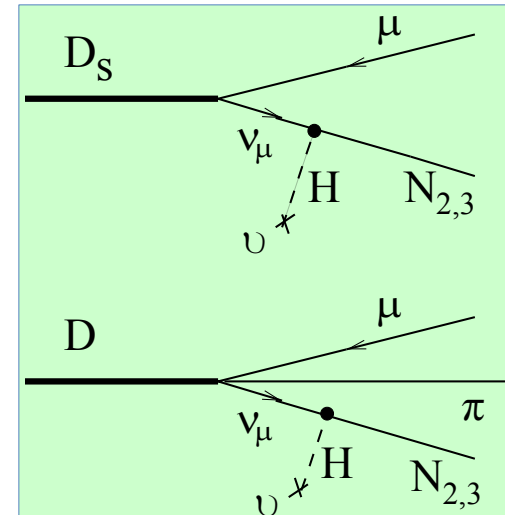
Production in charm decays...

... and decays to lighter SM particles

For what follows will focus on $N \rightarrow \mu^- \pi^+$ decay, $BF \sim 0.1-50\%$

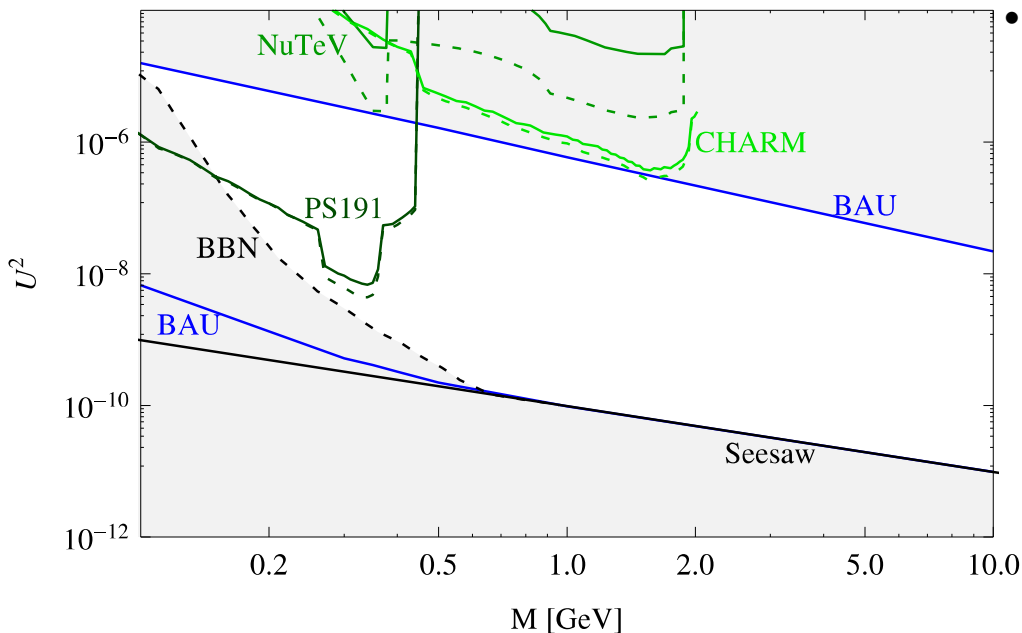
Similar BF for $N \rightarrow \mu^- \rho^+$,

$BF(N \rightarrow \nu \mu e) \sim 1-10\%$



Experimental and cosmological constraints

BAU, See-saw and Big Bang Nucleosynthesis (BBN) constraints indicate that previous experiments probed the interesting region only below the kaon mass :



• Previous searches :

- PS191('88)@PS 19.2 GeV
1.4×10¹⁹ pot, 128 m from target
- CHARM('86)@SPS 400 GeV, 2.4×10¹⁸ pot, 480 m from target
- NuTeV('99)@Fermilab 800 GeV, 2.5 × 10¹⁸ pot, 1.4 km from target

BBN, BAU and Seesaw give stronger constraints than experimental searches for $M_N > 400$ MeV

Experimental and cosmological constraints

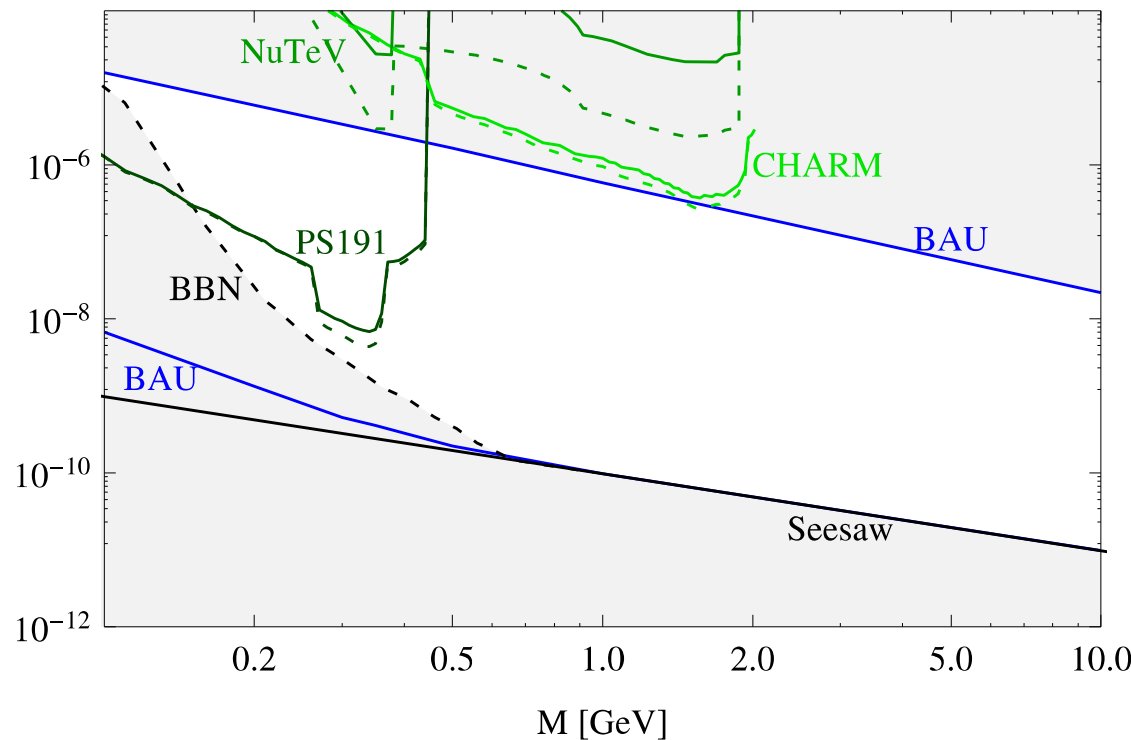
BAU, Seesaw and Big Bang Nucleosynthesis (BBN) constraints indicate that previous experiments probed the interesting region only below the kaon mass:

Mixing at both production and decay

BF($D \rightarrow NX$) around
 $10^{-8} - 10^{-12}$

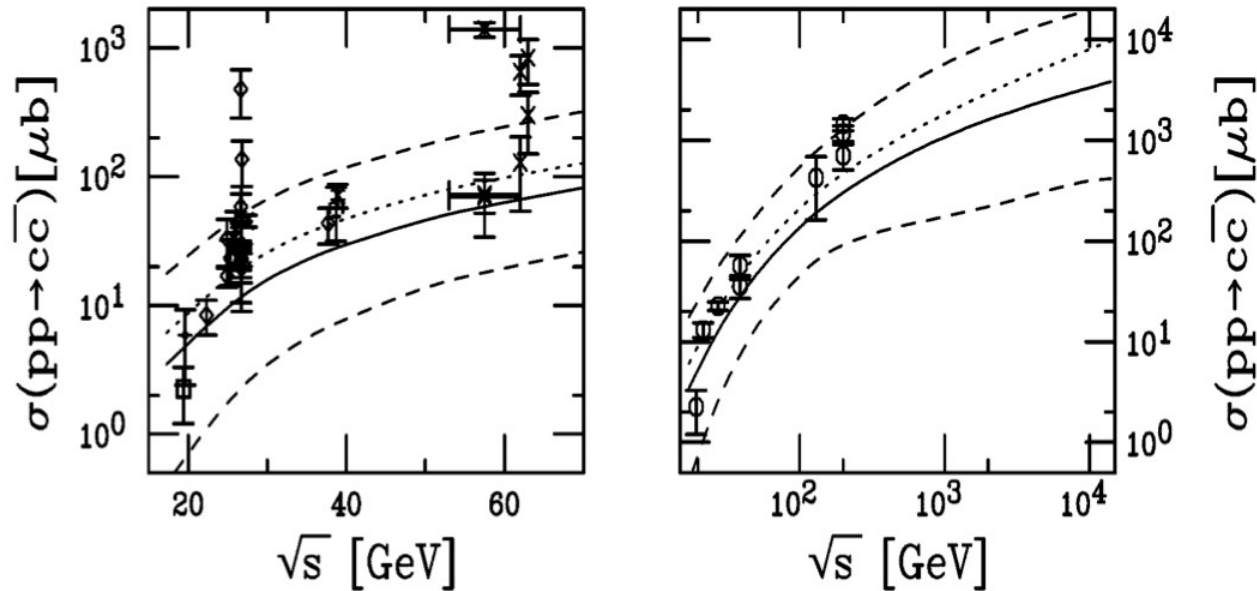
Lifetime can give
further factor 10^{-4}

Need $>10^{16}$ D mesons!



Where to produce charm?

$c\bar{c}$ cross-section:



arxiv.org/pdf/0709.2531v1

LHC ($\sqrt{s} = 14$ TeV)

1 ab^{-1} (i.e. 3-4 years), $\sim 2 \times 10^{16}$ in 4π

SPS (400 GeV p-on-target (pot) $\sqrt{s} = 27$ GeV)

2×10^{20} pot (i.e. 3-4 years): $\sim 2 \times 10^{17}$

Fermilab: 120 GeV, $10 \times$ smaller $\sigma_{c\bar{c}}$, $10 \times$ pot by 2025 for LBNE

Could mass range be extended?

Would neutrinos from B-decays extend the mass range for $N_{2,3}$ upwards?

Produced with factor 20-100 smaller cross-section

Dominant semi-leptonic decay $B \rightarrow D\mu\nu$

Similar limit for neutrino mass

Charmless $B \rightarrow \pi\mu\nu$ heavily Cabibbo suppressed

B decays are not at all competitive

Decays from $Z \rightarrow \nu N$ neither competitive

Lifetime too long to contain

Experimental design

Propose a beam dump experiment at the CERN SPS with a total of $\sim 2 \times 10^{20}$ protons on target

Crucial design parameters:

- Minimise residual neutrino and muon fluxes

 - Can produce K^0 that decay in detector and mimic signal events

- Short-lived resonances generate 10^9 muons/spill

 - Muon shield

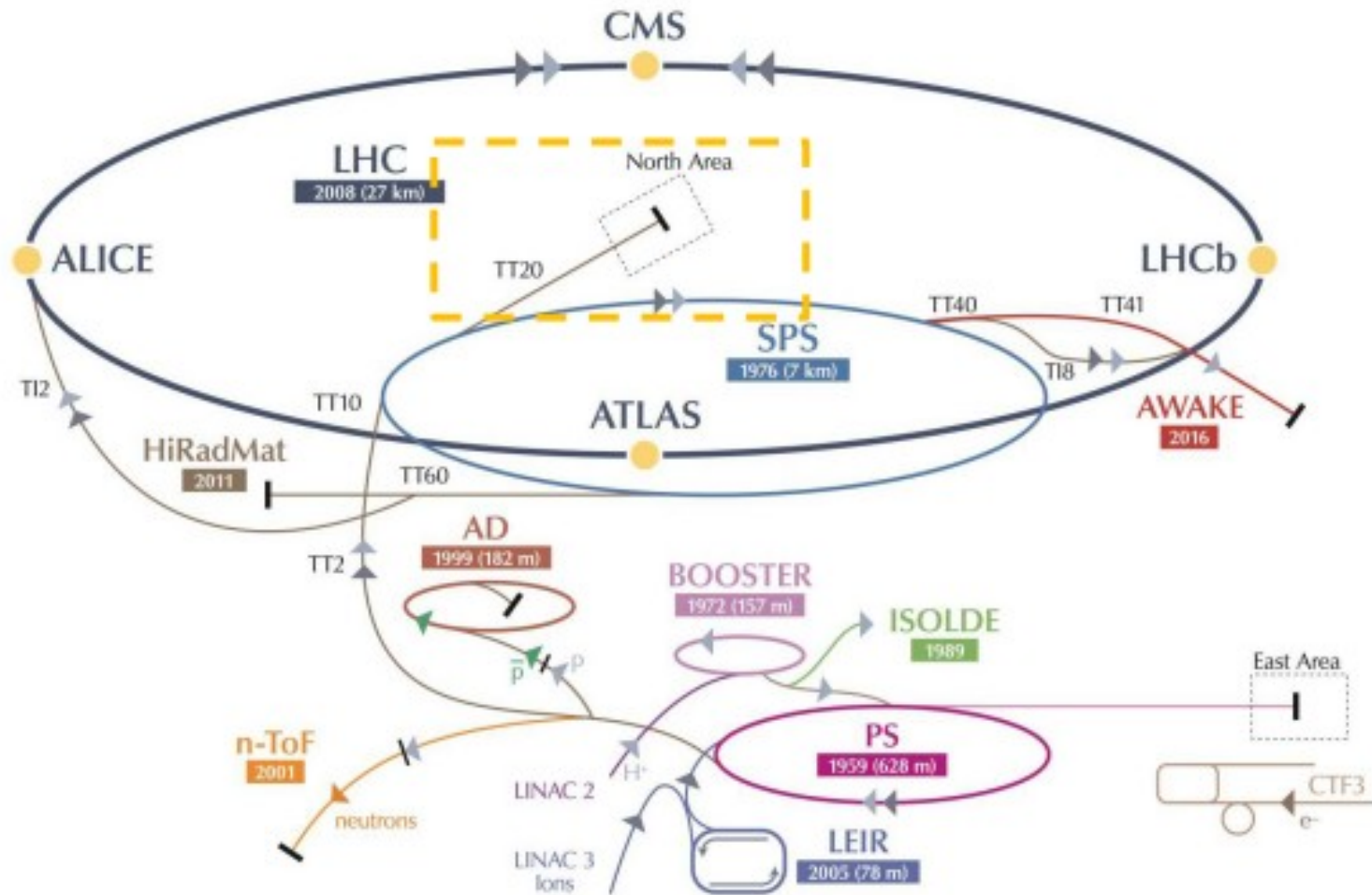
- Neutrinos from light meson decays

 - Dense target/hadron absorber

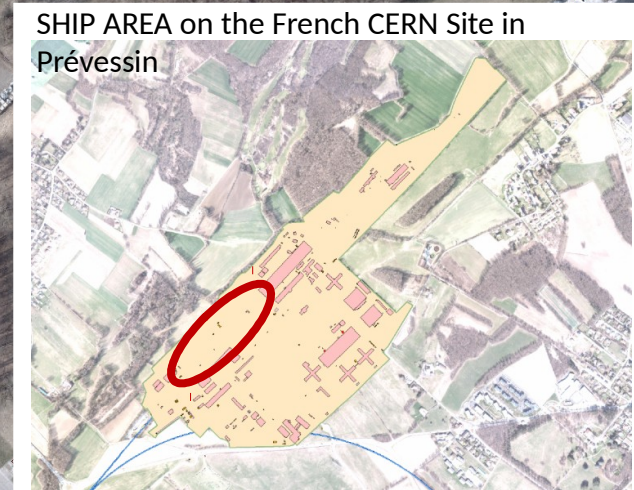
- Prevent neutrino interactions from mimicking HNL decay

 - Evacuate decay volume

Beamline



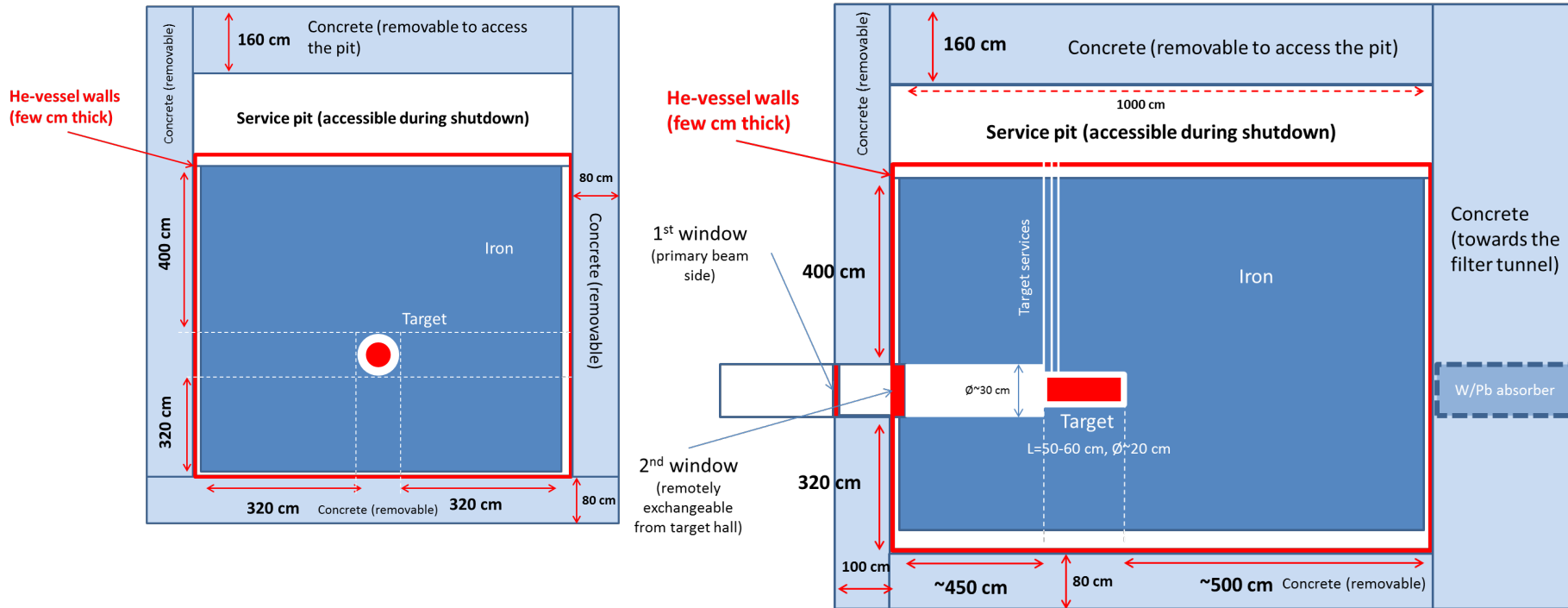
Target



Target

Compact tungsten target to minimise neutrinos from kaon and pion decays

Significant (but achievable) requirement on cooling and radiation protection.



Timeline for civil engineering

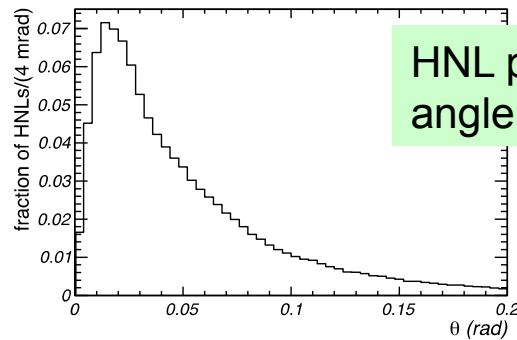
The LS2 of of the LHC is critical for building the new extraction point from the SPS

This is what drives the aggressive schedule for a Technical Proposal

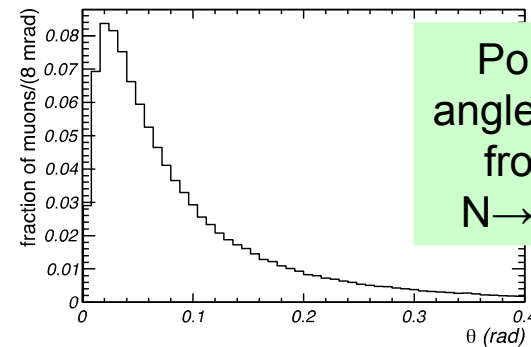
SHIP_CE DRAFT PLANNING	2014				2015				2016				2017				2018				2019				2020				2021				2022			
	Q1	Q2	Q3	Q4	Q1	Q2	Q3	Q4	Q1	Q2	Q3	Q4	Q1	Q2	Q3	Q4	Q1	Q2	Q3	Q4	Q1	Q2	Q3	Q4	Q1	Q2	Q3	Q4	Q1	Q2	Q3	Q4	Q1	Q2	Q3	Q4
LHC operation	Red				Green				Green				Green				Red				Yellow				Green				Green							
SPS operation	Red				Green				Green				Green				Red				Yellow				Green				Green							
Technical Proposal	Blue				Blue				Blue				Blue				Blue				Blue				Blue				Blue							
SHIP Project approval	Blue				Blue				Blue				Blue				Blue				Blue				Blue				Blue							
Pre-construction activities(Design, tendering, permits)	Blue				Blue				Blue				Blue				Blue				Blue				Blue				Blue							
CE works for extraction tunnel, target area	Blue				Blue				Blue				Blue				Blue				Blue				Blue				Blue							
CE works for TDC2 junction cavern	Blue				Blue				Blue				Blue				Blue				Blue				Blue				Blue							
CE works for filter tunnel and detector hall	Blue				Blue				Blue				Blue				Blue				Blue				Blue				Blue							
GS-SE Resources	Orange				Orange				Orange				Orange				Orange				Orange				Orange				Orange							
Engineer	Orange				Orange				Orange				Orange				Orange				Orange				Orange				Orange							
Technician/Fellow	Green				Green				Green				Green				Green				Green				Green				Green							
Draughtman	Red				Red				Red				Red				Red				Red				Red				Red							

Experimental design

HNLs produced in charm decays have significant p_T



HNL polar
angle



Polar
angle of μ
from
 $N \rightarrow \mu \pi$

Detector must be close to target to maximise geometrical acceptance

Shielding for muons must be as short as possible

Secondary beam line

Initial reduction of beam induced backgrounds

- Heavy target (50 cm of W)
- Hadron absorber
- Muon shield: optimization of active and passive shields is underway

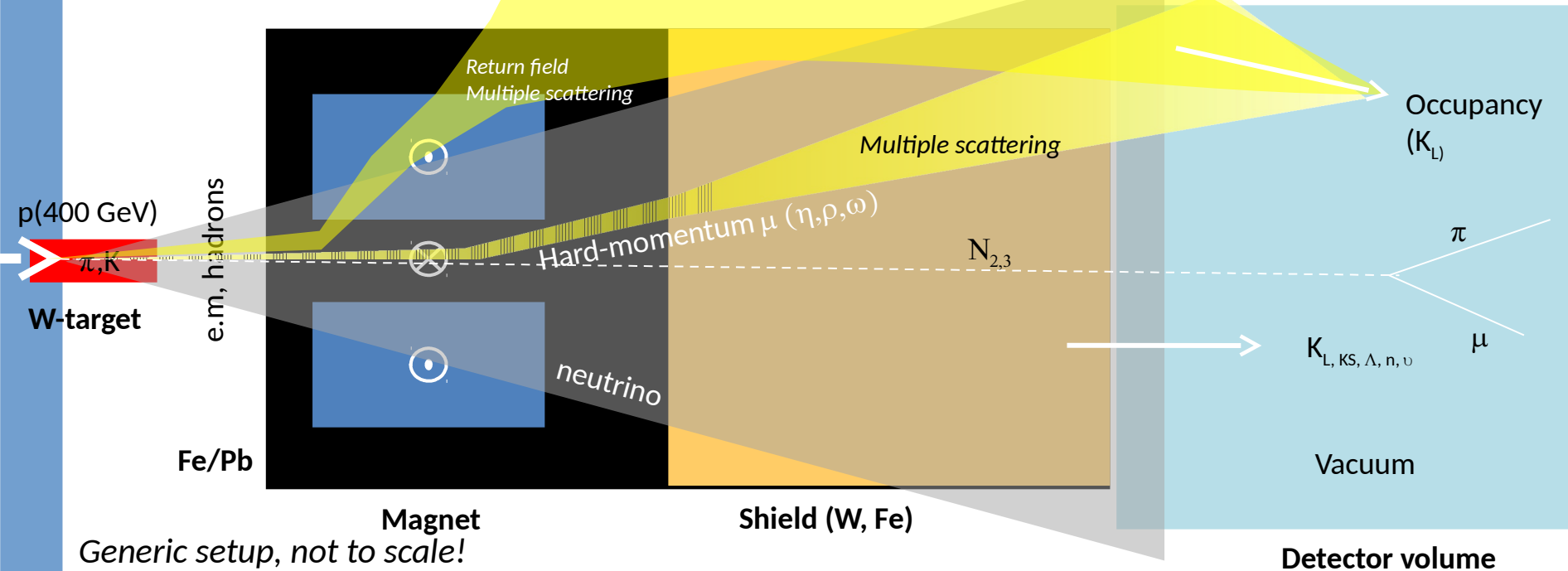
Acceptable occupancy < 1% per spill of 5×10^{13} p.o.t.

spill duration 1s $\Rightarrow < 50 \times 10^6$ muons

spill duration 10ms $\Rightarrow < 50 \times 10^3$ muons

spill duration 10 μ s $\Rightarrow < 500$ muons

Low-mid-momentum μ from fast decays of π, K



Muon Shield

Without μ -filter: 5×10^9 / spill (5×10^{13} pot)

Idea to reduce background from μ -interactions to below ν -background

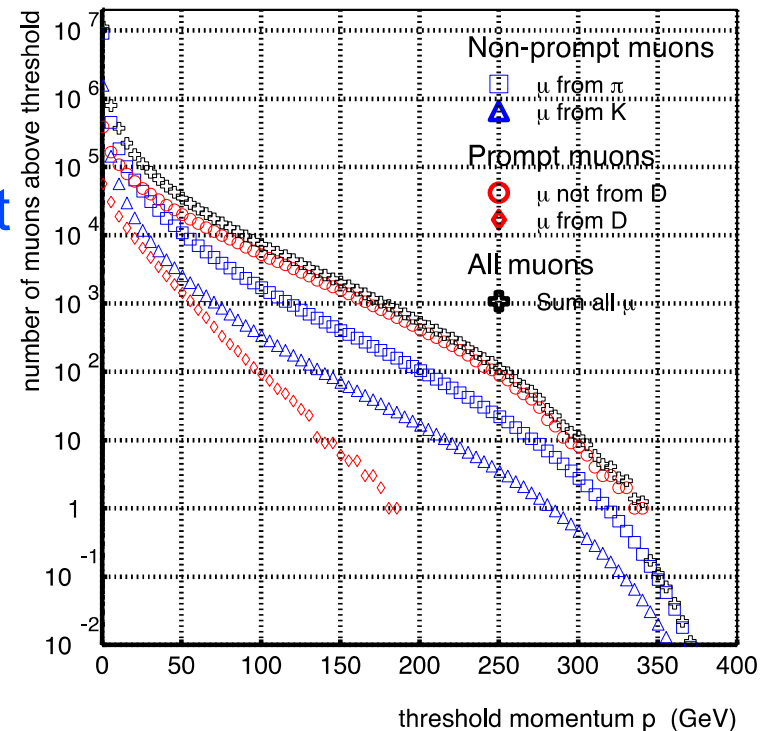
Acceptable rate $\sim 10^5 \mu / 2 \times 10^{20}$ pot

Main sources of muons simulated using PYTHIA

Two alternatives for shield:

Passive: i.e. use high Z material: need 54 m of W to stop 400 GeV μ

Active (+passive): need 40 Tm to deflect 400 GeV μ outside acceptance



Neutrino backgrounds

Neutrino interactions in the decay volume :

After shield expect 2×10^4 per 2×10^{20} pot at atmospheric pressure

Negligible at 0.01 mbar

Neutrino interactions in the final part of the muon shield :

Use GEANT and GENIE to simulate the CC and NC neutrino interactions

CC(NC) rate of $\sim 6(2) \times 10^5$ per interaction length per 2×10^{20} pot

Use veto-station to suppress short lived

$\nu_\mu + p \rightarrow X + K_L \rightarrow \mu\pi\nu$ main background

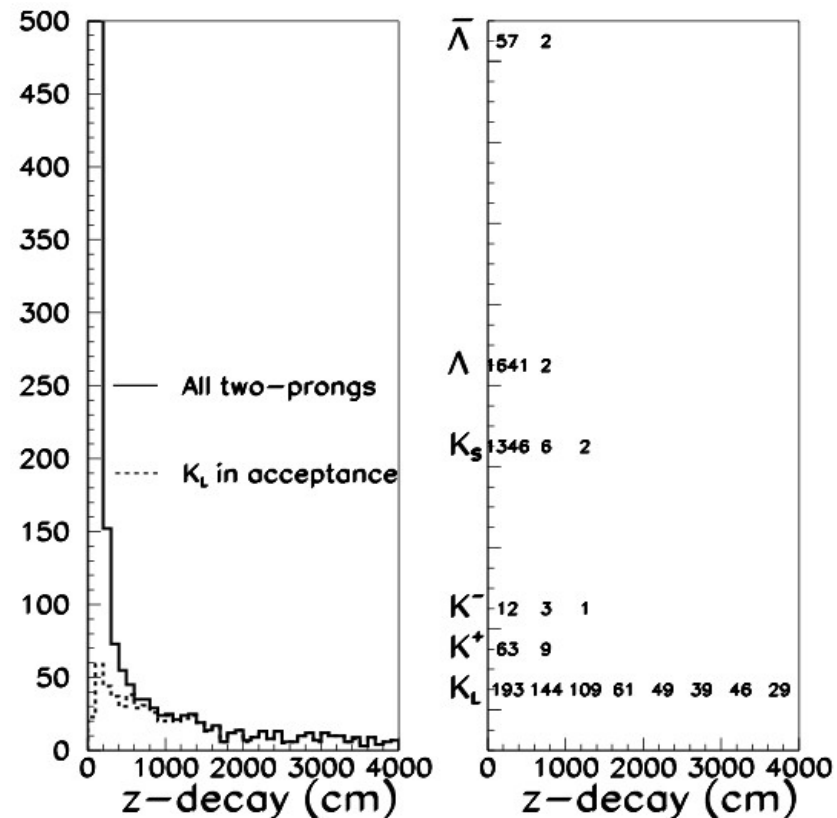
Requiring μ -id. for one of the two decay products

→ 150 two-prong vertices in 2×10^{20} pot

Neutrino backgrounds

Neutrino interactions in the decay volume

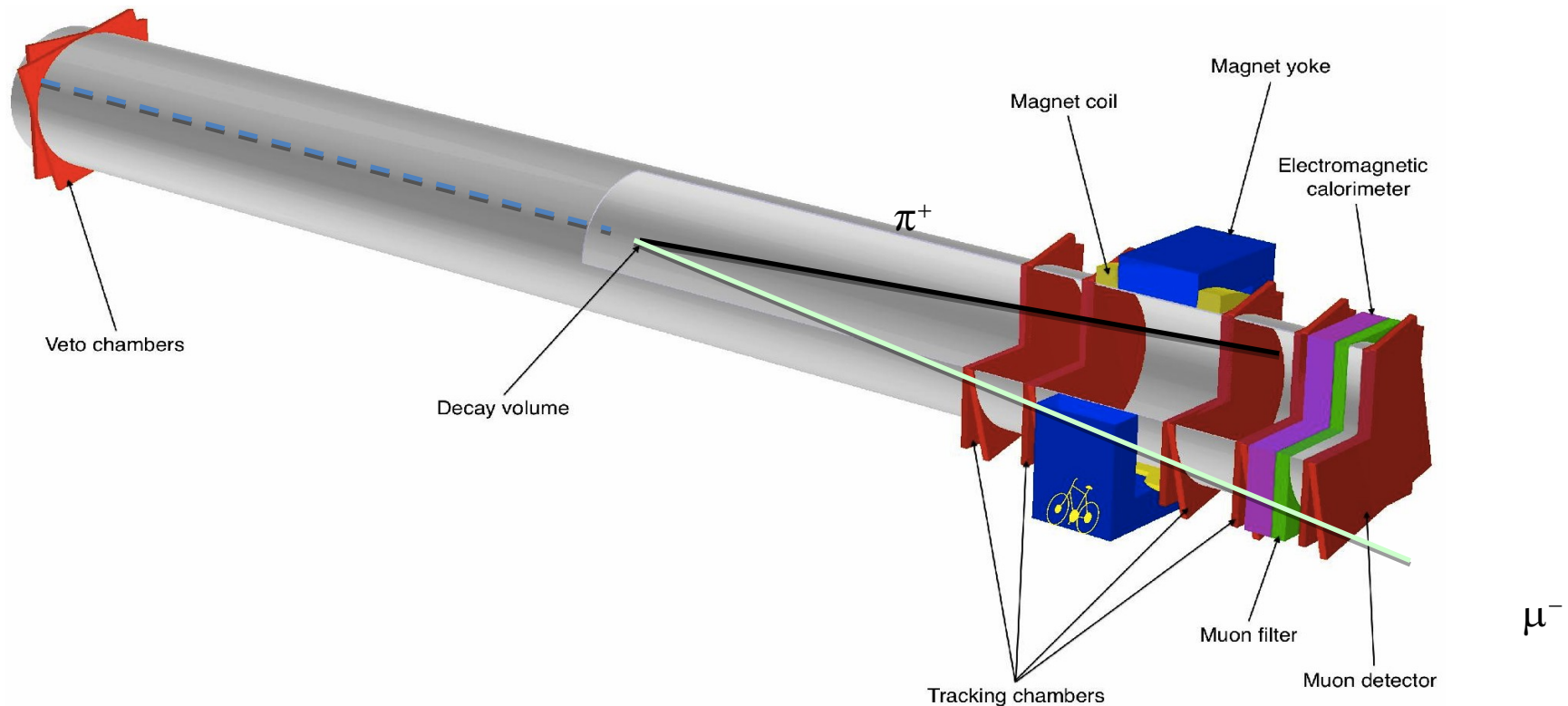
~10% of neutrino interactions in the muon shield just upstream of the decay volume produce Λ or K^0



Detector Concept

Aim to reconstruct HNL decays into the final states: $\mu^- \pi^+$, $\mu^- \rho^+$, $e^- \rho^+$

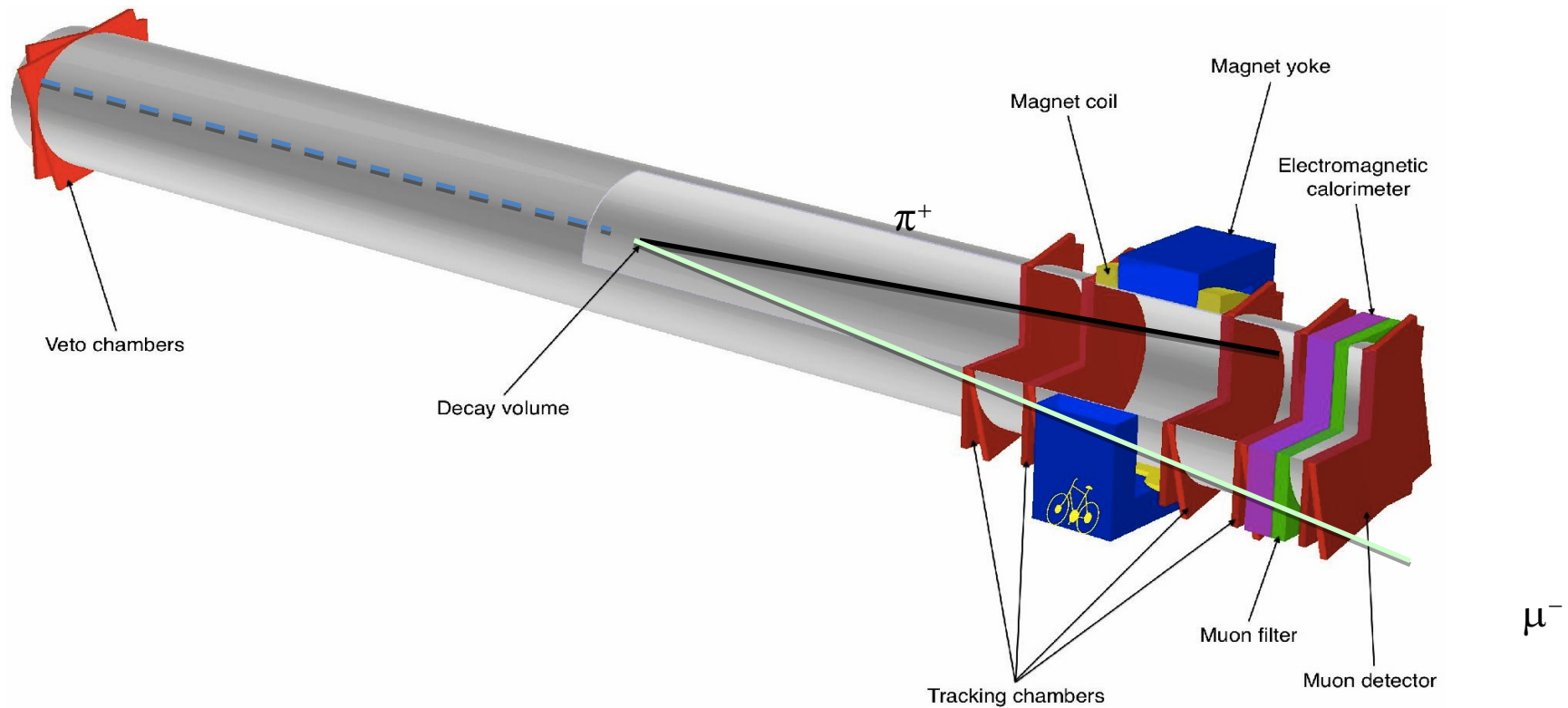
Require long decay volume, magnetic spectrometer, muon detector and electromagnetic calorimeter



Detector Concept

5 m diameter, 50 m length vacuum vessel

10 m long magnetic spectrometer with 0.5 Tm dipole magnet and four tracking chambers



Detector Technologies

Dipole magnet

Magnet similar to LHCb design required,
but with $\sim 40\%$ less iron and $3\times$ less power

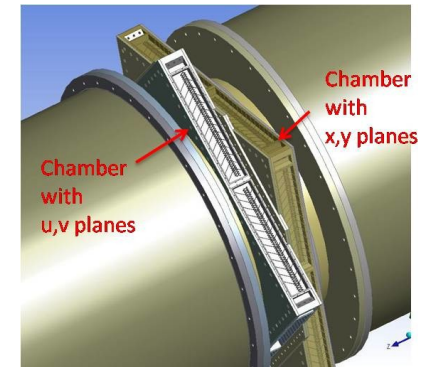
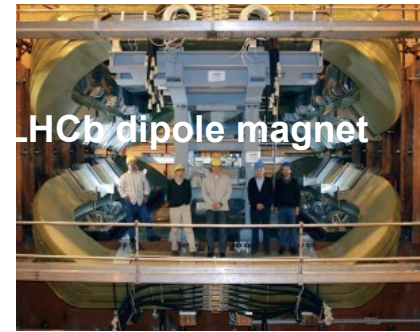
Free aperture of $\sim 16 \text{ m}^2$

Field integral $\sim 0.5 \text{ Tm}$ over 5 m length

Vacuum tank and straw tracker

NA62 has 10^{-5} mbar pressure, only 10^{-2} mbar here

Have demonstrated gas tightness of straw tubes with $120 \mu\text{m}$ spatial resolution and 0.5% X^0 material budget in long term tests



Detector Technologies

Electromagnetic calorimeter

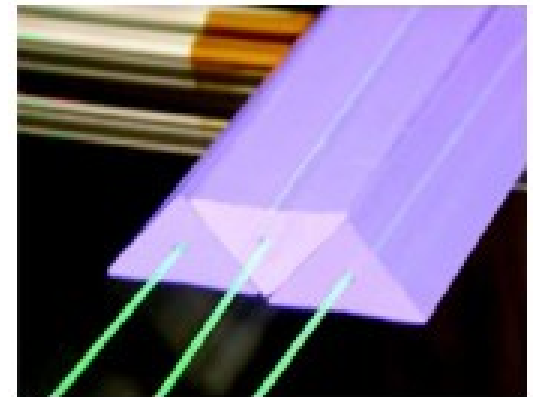
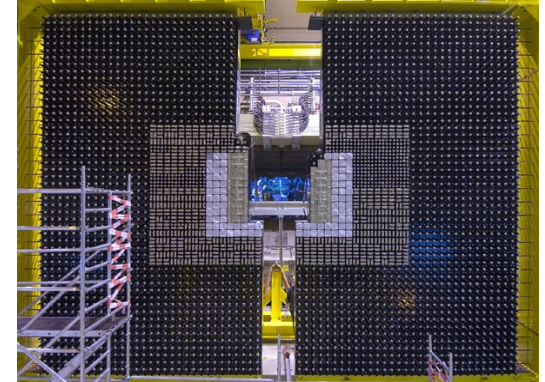
Shashlik technology used in LHCb would provide economical solution with good energy and time resolution

Muon detector

Scintillator strips with WLS fibres and Silicon Photomultiplier (SiPM) an attractive option

Trigger and DAQ

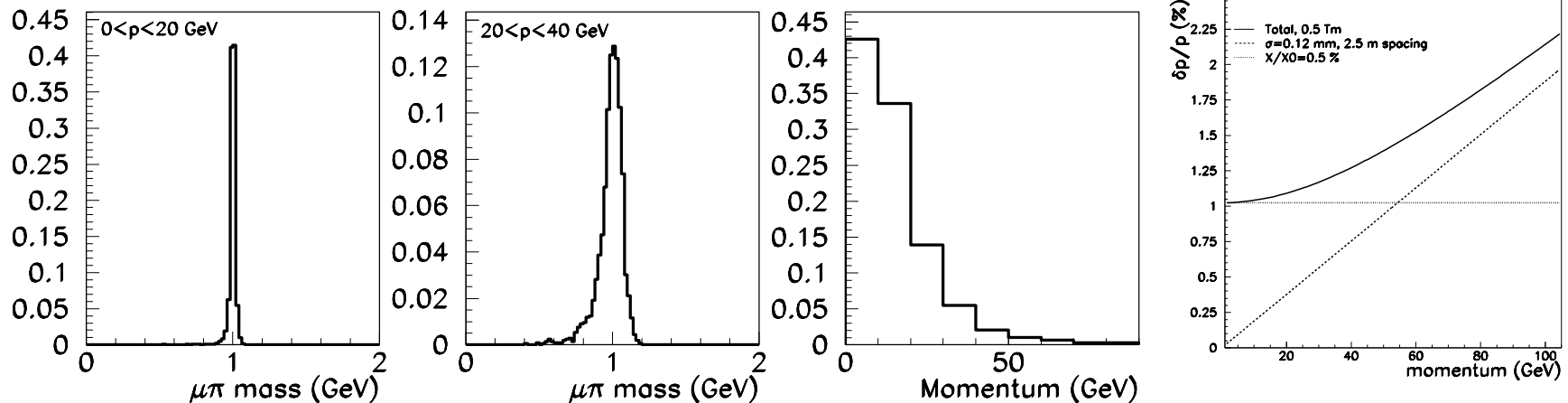
Requirements on both are very modest due to low data rate



Spectrometer resolution

Arrange spectrometer such that multiple scattering and spatial resolution of straw tubes give similar contribution to the overall $\Delta p/p$

For $m(N_{2,3}) = 1 \text{ GeV}$, 75% of $\mu\pi^+$ decay products have $p < 20 \text{ GeV}$



For 0.5 Tm field integral $\sigma(\text{mass}) \sim 40 \text{ MeV}$ for $p < 20 \text{ GeV}$

Good discrimination between high mass tail from small number of residual $K_L \rightarrow \mu\pi^0$ and 1 GeV HNL

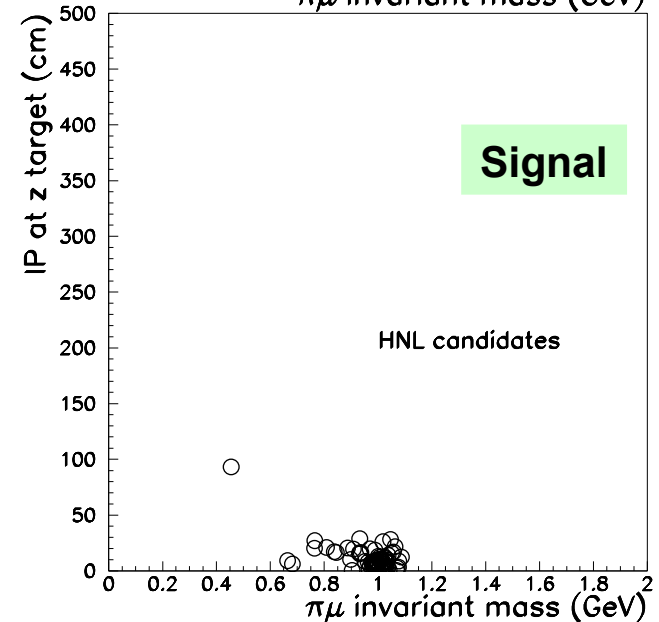
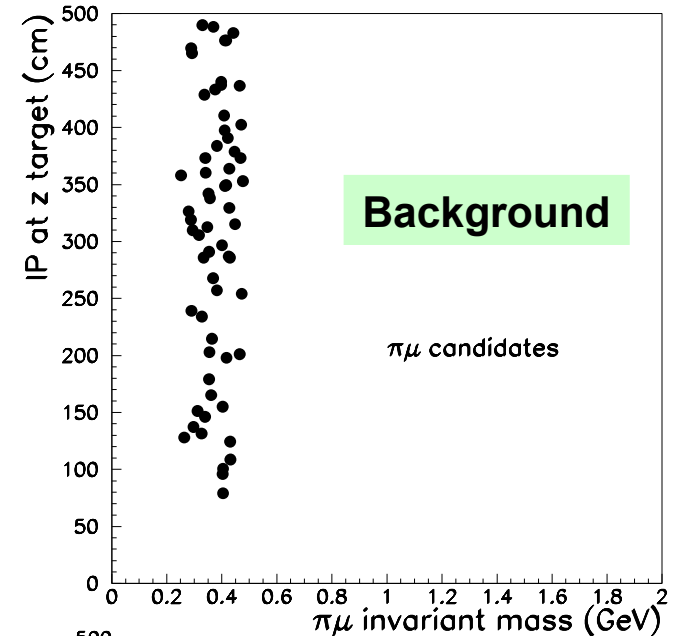
Residual Backgrounds

K_L produced in the final part of the muon shield have very different pointing to the target compared to signal events

Use Impact Parameter (IP) to further suppress K_L background

IP < 1 m is 100% efficient for signal and leaves only a handful of background events

The IP cut will also be used to reject backgrounds induced in neutrino interactions in the material surrounding the detector



Expected event yield

Integral mixing angle U^2 is given by $U^2 = U_e^2 + U_\mu^2 + U_\tau^2$

Make a conservative estimate of the sensitivity by only considering the decay $N_{2,3} \rightarrow \mu^- \pi^+$ – probes U_μ^2

Expected number of signal events then,

$$N_{\text{signal}} = n^{\text{pot}} \times 2\chi_{\text{cc}} \times \text{BR}(U_\mu^2) \times \varepsilon_{\text{det}}(U_\mu^2)$$

Strongest experimental limit for $M_N \sim 1$ GeV at $U_\mu^2 = 10^{-7}$

Would then expect $\tau_N = 1.8 \times 10^{-5}$ s and ~ 12 k fully reconstructed $N \rightarrow \mu^- \pi^+$

For cosmologically favoured region $U_\mu^2 = 10^{-8}$ ($\tau_N = 1.8 \times 10^{-4}$ s)

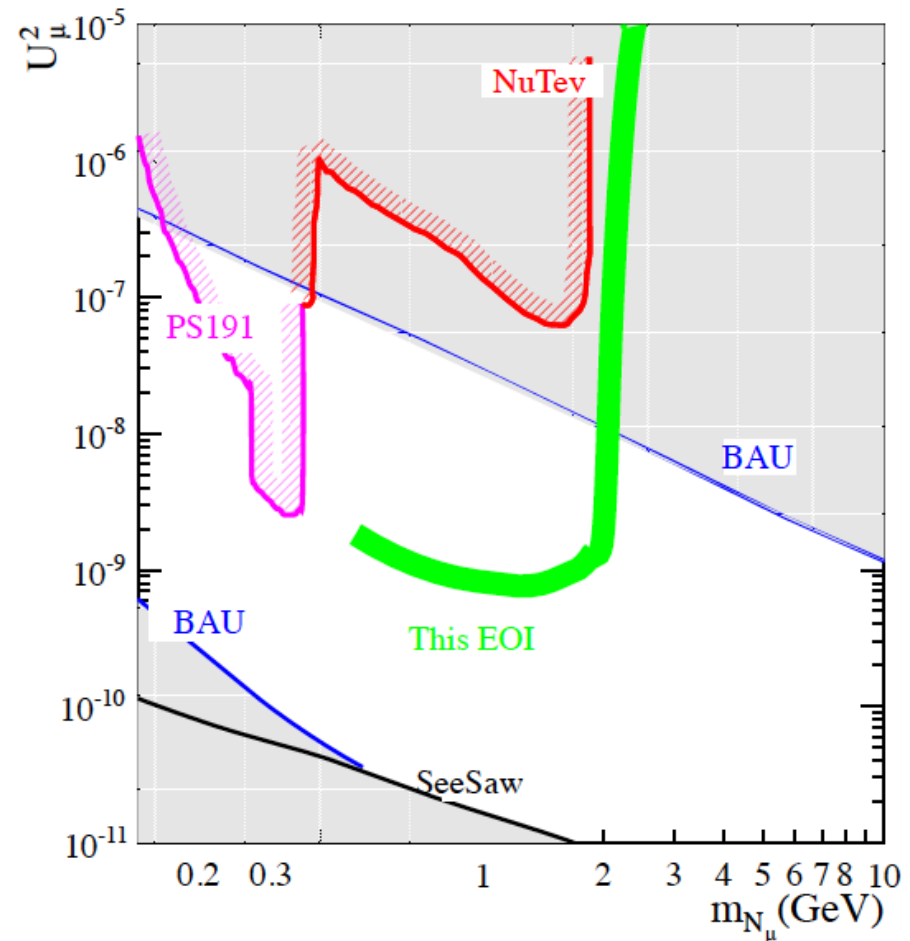
Would expect 120 fully reconstructed events

Expected sensitivity

For $M_N < 2$ GeV the proposed experiment has discovery potential for the cosmologically favoured region with $10^{-7} < U_\mu^2 < \text{a few} \times 10^{-9}$

Limit from decay channels with electromagnetic not studied yet.

Will extend search on U^2 and limit on U_e^2



Status of the SPSC review

Submitted our EOI in Oct 2013 [CERN-SPSC-2013-024 / SPSC-EOI-010 / arXiv:1310.1762]

SPSC assigned four referees – provided answers to their questions [http://ship.web.cern.ch/ship/EOI/SPSC-EOI-010_ResponseToReferees.pdf]

SPSC discussed our proposal Jan 2015, official feedback

"The Committee **received with interest** the response of the proponents to the questions raised in its review of EOI010.

The SPSC **recognises** the interesting physics potential of searching for heavy neutral leptons and investigating the properties of neutrinos.

Considering the large cost and complexity of the required beam infrastructure as well as the significant associated beam intensity, such a project should be designed as a general purpose beam dump facility with the broadest possible physics programme, including maximum reach in the investigation of the hidden sector.

To further review the project the Committee **would need** an extended proposal with further developed physics goals, a more detailed technical design and a stronger collaboration."

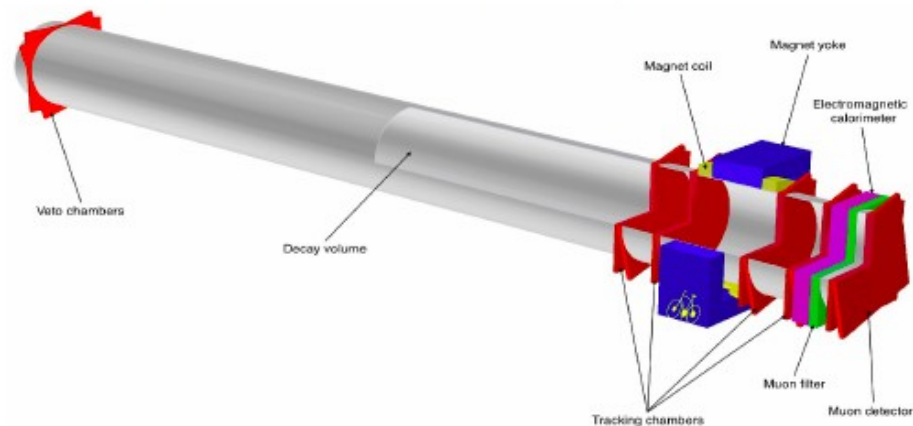
A strengthening collaboration



SHIP - Search for Hidden Particles

CERN, Universität Zürich, EPFL Lausanne, INFN Cagliari, Università Federico II and INFN Napoli, Imperial College London

Experiment to search for Heavy Neutral Leptons at the SPS



We propose a new fixed-target experiment at the CERN SPS accelerator to search for *hidden particles*. In particular, to search for Heavy Neutral Leptons (HNLs) produced in charm decays. HNLs are right-handed partners of the Standard Model neutrinos. The existence of such particles is strongly motivated by theory, as they can simultaneously explain the baryon asymmetry of the Universe, account for the pattern of neutrino masses and oscillations, and provide a Dark Matter candidate.

SHIP is a collaboration of six institutes: CERN, Universität Zürich, École Polytechnique Fédérale de Lausanne, INFN Sezione di Cagliari, Università Federico II and INFN Napoli, Imperial College London. Groups interested in joining should contact [Andrey Golutin](#) and [Jaap Panman](#). The extension of the collaboration will be discussed at the [First SHIP Workshop](#) that will take place in **Zürich the 10-12 June 2014**.

First SHiP Workshop, 10-12th June 2014

Around 100 people met in Zurich and discussed

Physics reach of SHiP detector

Detector requirements and technologies

A proto collaboration is now in place ...

Theoretical Overview (10th June)

Review of heavy neutral leptons, with discussions about leptogenesis and cosmological constraints

Theory review (11th June Morning)

Discussion of theoretical status and present experimental constraints

Facility and Experiment (11th June Afternoon)

Discussion on the primary beam line, target and detector design for the SHiP experiment

Tau neutrinos and SHiP detector (12th June Morning)

Discussion on the electronics and DAQ system for the SHiP experiment and on the detector for tau neutrinos

Summary and discussion (12th June Afternoon)

Conclusions

The proposed experiment will search for NP in the largely unexplored domain of new, very weakly interacting particles with masses below the Fermi scale

Detector is based on existing technologies

Ongoing discussion of the beam line with CERN experts

The discovery of a HNL would have enormous impact – could solve several of the significant problems of the SM

- The origin of the baryon asymmetry of the Universe

- The origin of neutrino mass

- The results of this experiment, together with cosmological and astrophysical data, could be crucial to determine the nature of Dark Matter

Wide range of other hidden sector physics under investigation